

CGC and Initial Conditions for Heavy Ion Collisions

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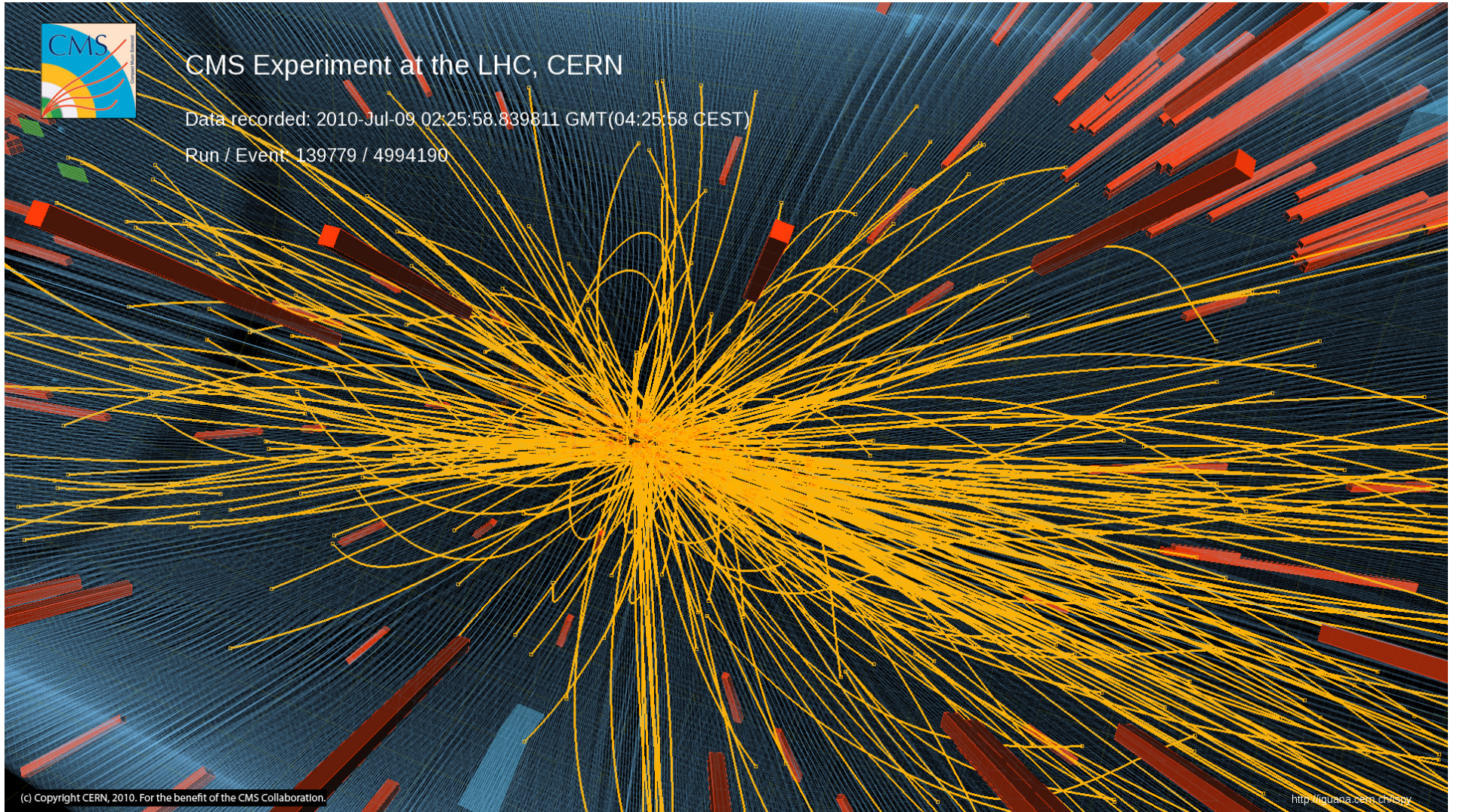
August 22nd, 2012

NC STATE UNIVERSITY

Two particle correlations in p+p collisions

1. Evidence for saturation dynamics (the p+p near-side Ridge)
2. Evidence for BFKL evolution (mini-jet decorrelation)
3. Prospects for p+Pb @ LHC

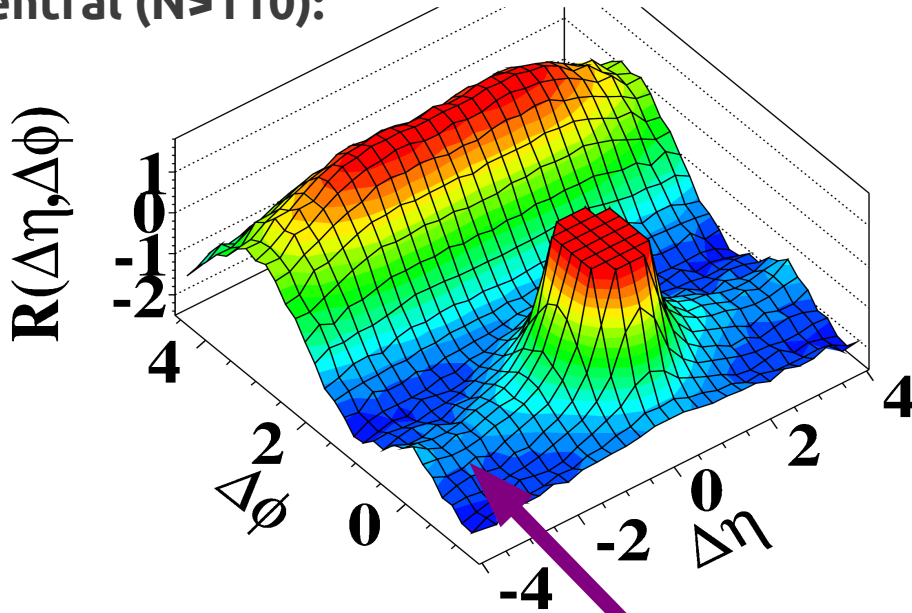
Part 1: the p+p near side ridge



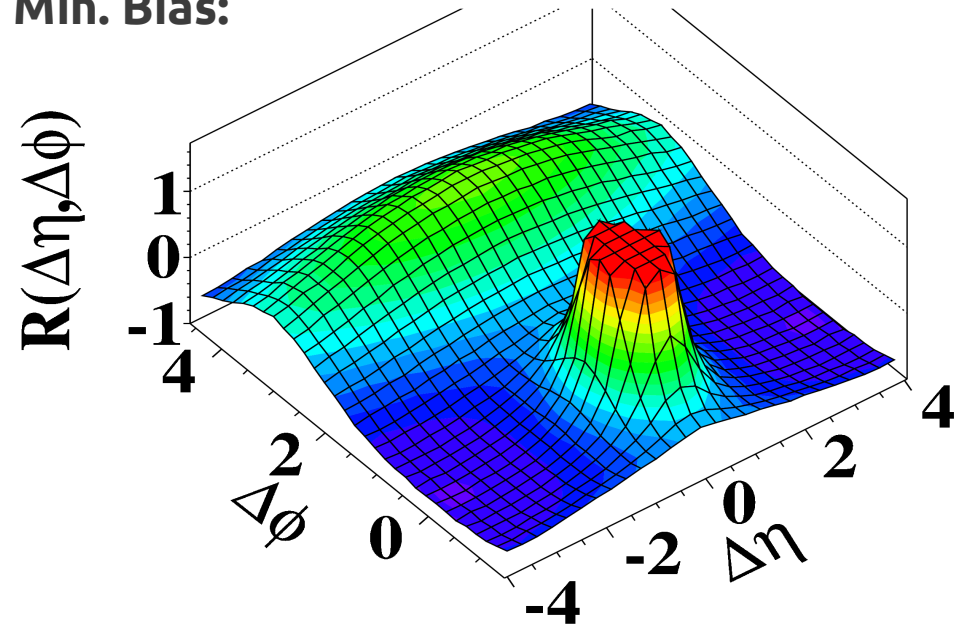
More than 200 charged particles!

$$1.0 < p_T [\text{GeV}] < 3.0$$

(d) CMS $N \geq 110$, $1.0 \text{ GeV}/c < p_T < 3.0 \text{ GeV}/c$
Central ($N > 110$):

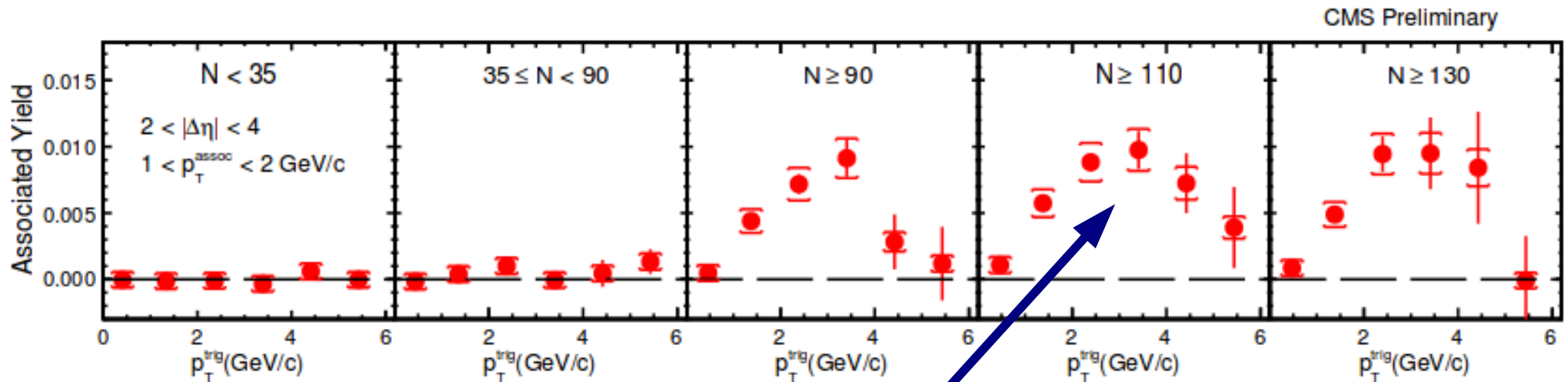


(b) CMS MinBias, $1.0 \text{ GeV}/c < p_T < 3.0 \text{ GeV}/c$
Min. Bias:



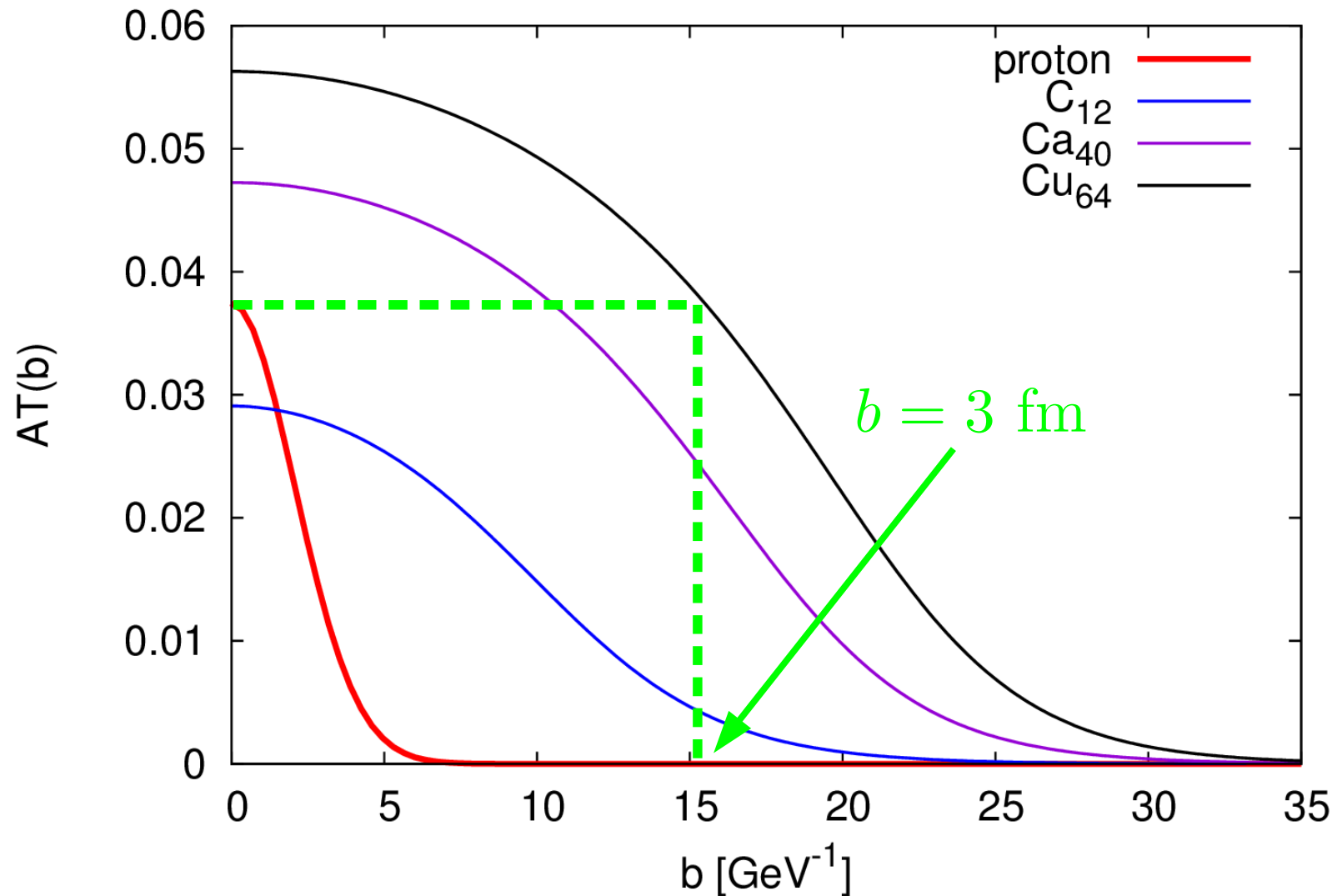
Novel structure seen in highest multiplicity events

The Ridge



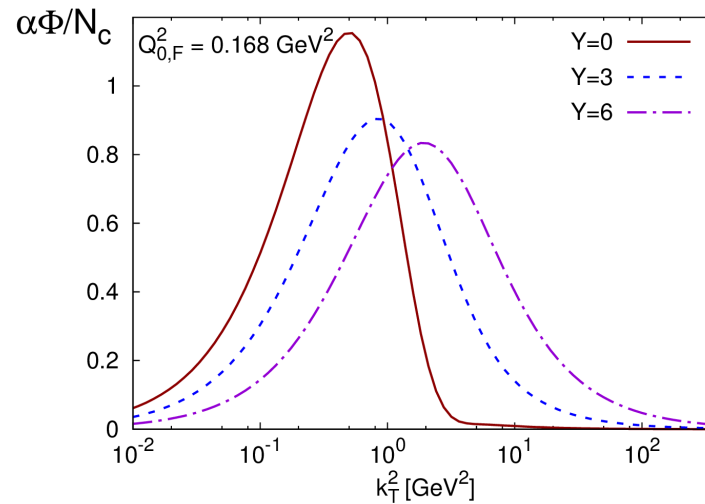
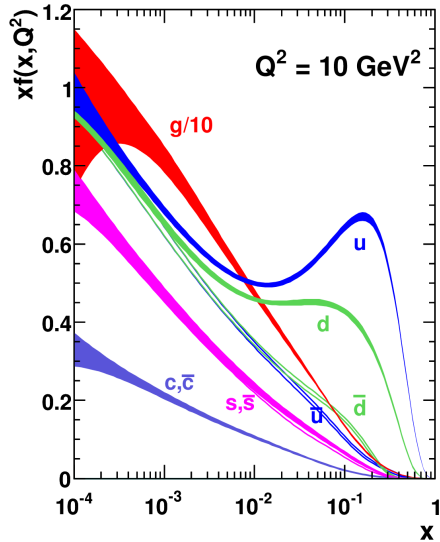
1. There is a clear scale in the data
2. It is semi-hard and will be argued to be related to Q_s

Multiplicity the same as in Cu+Cu !



The proton pre-collision

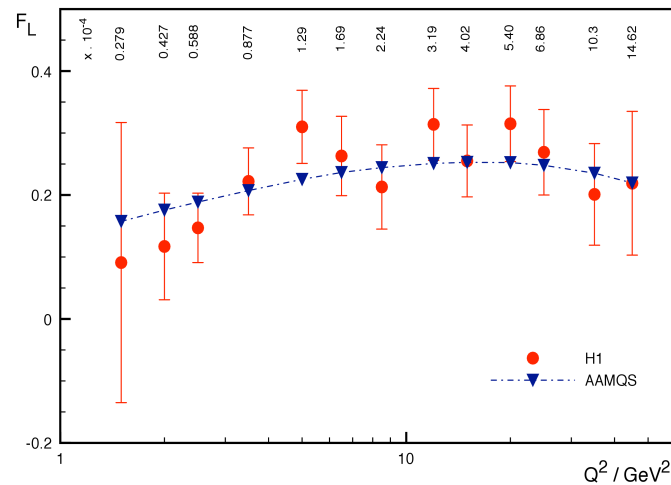
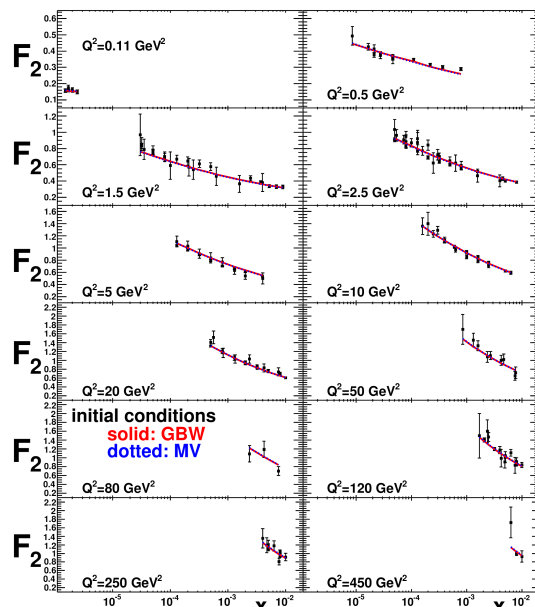
Our field has a good understanding of the proton wave-function:



NLO DGLAP fits:
<http://mstwpdf.hepforge.org/>

NLO-BK:
 Balitsky, Chirilli PRD 77 014019
 Kovchegov, Weigert NPA 784 188
 Albacete, Kovchegov PRD 75 125021

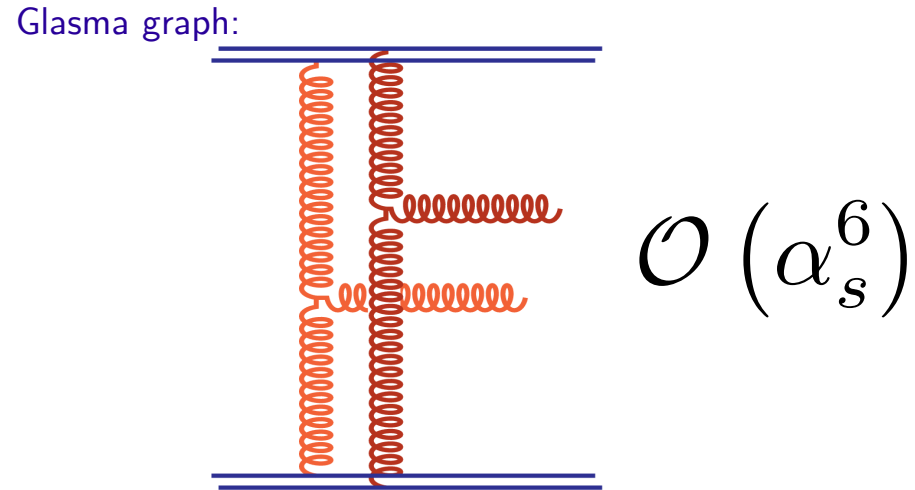
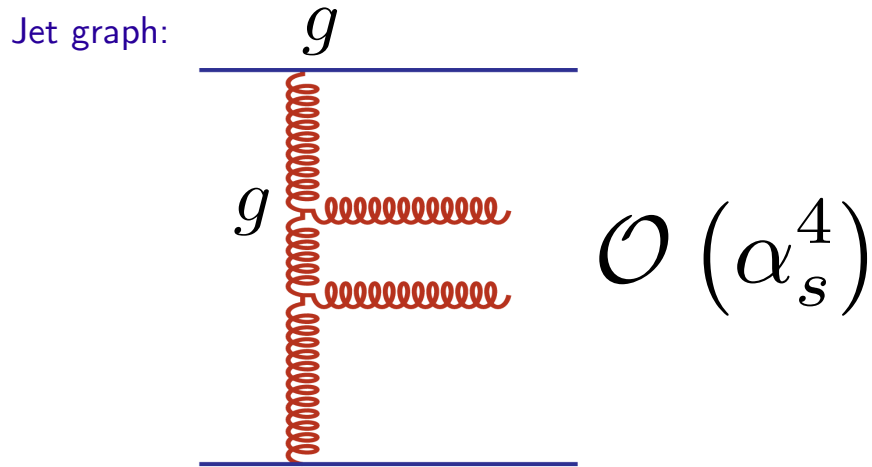
15 years of HERA data support this picture:



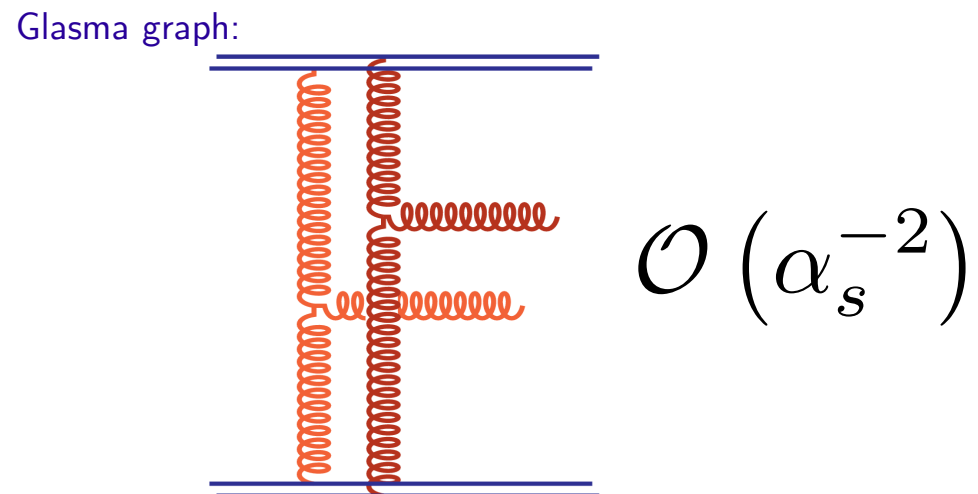
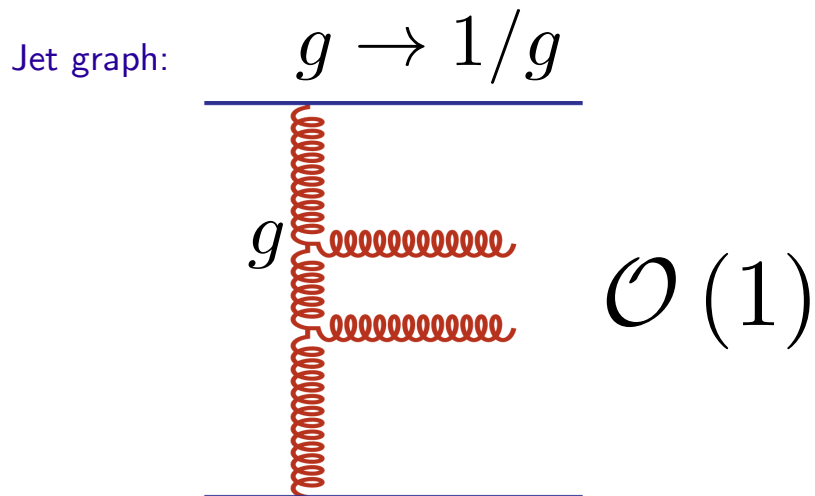
Albacete, Milhano, Quiroga-Arias, Rojo, arXiv:1203.1043 (2012).
 Quiroga-Arias, Albacete, Armesto, Milhano, Salgado, J.Phys.G G38 (2011) 124124.
 Albacete, Armesto, Milhano, Salgado, PRD80 (2009) 034031.

Power counting in QCD: multiparticle production

Low color charge density (min bias):

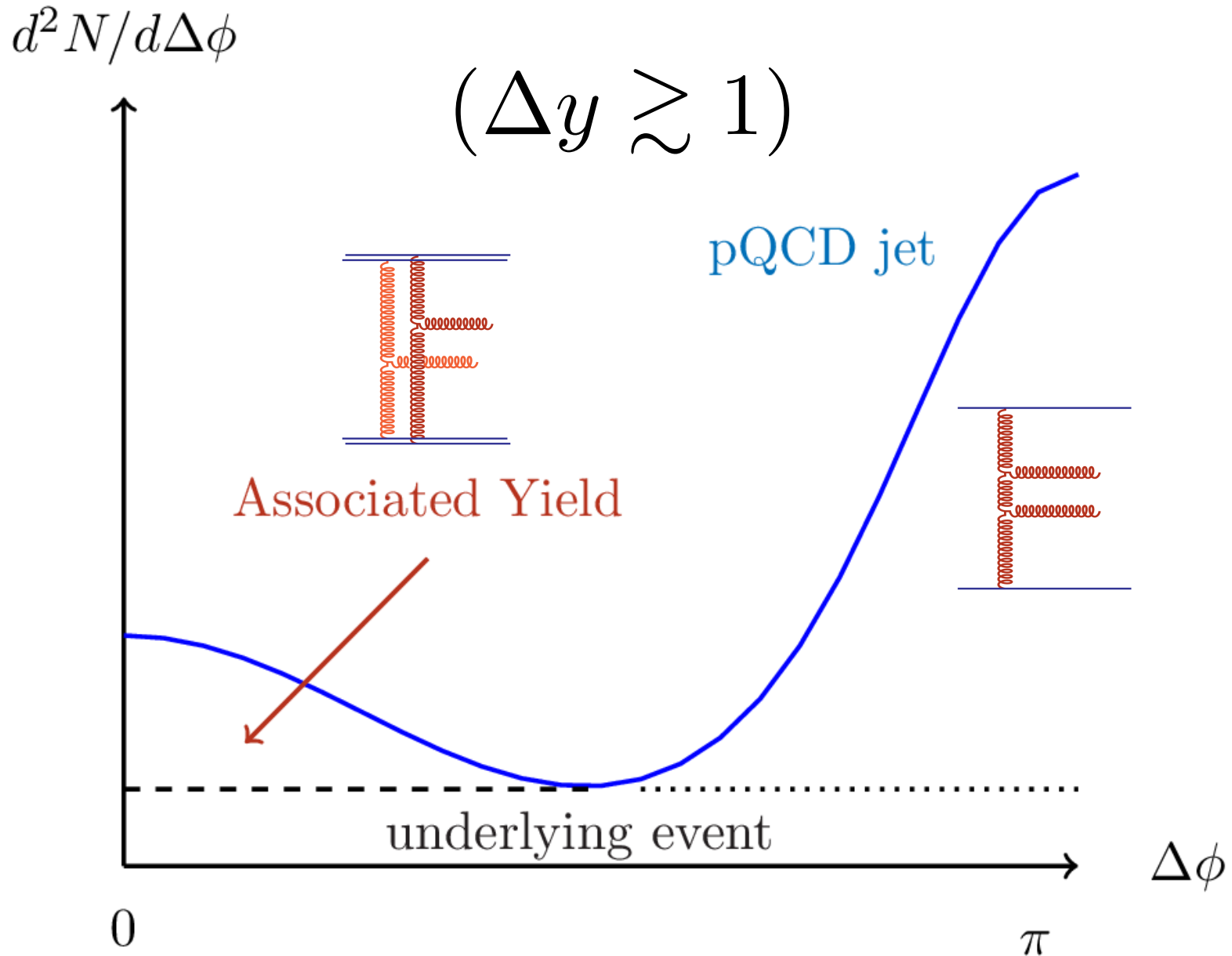


High color charge density (central):



Expect α_s^8 enhancement of “Glasma” graph! Is this seen in the data?

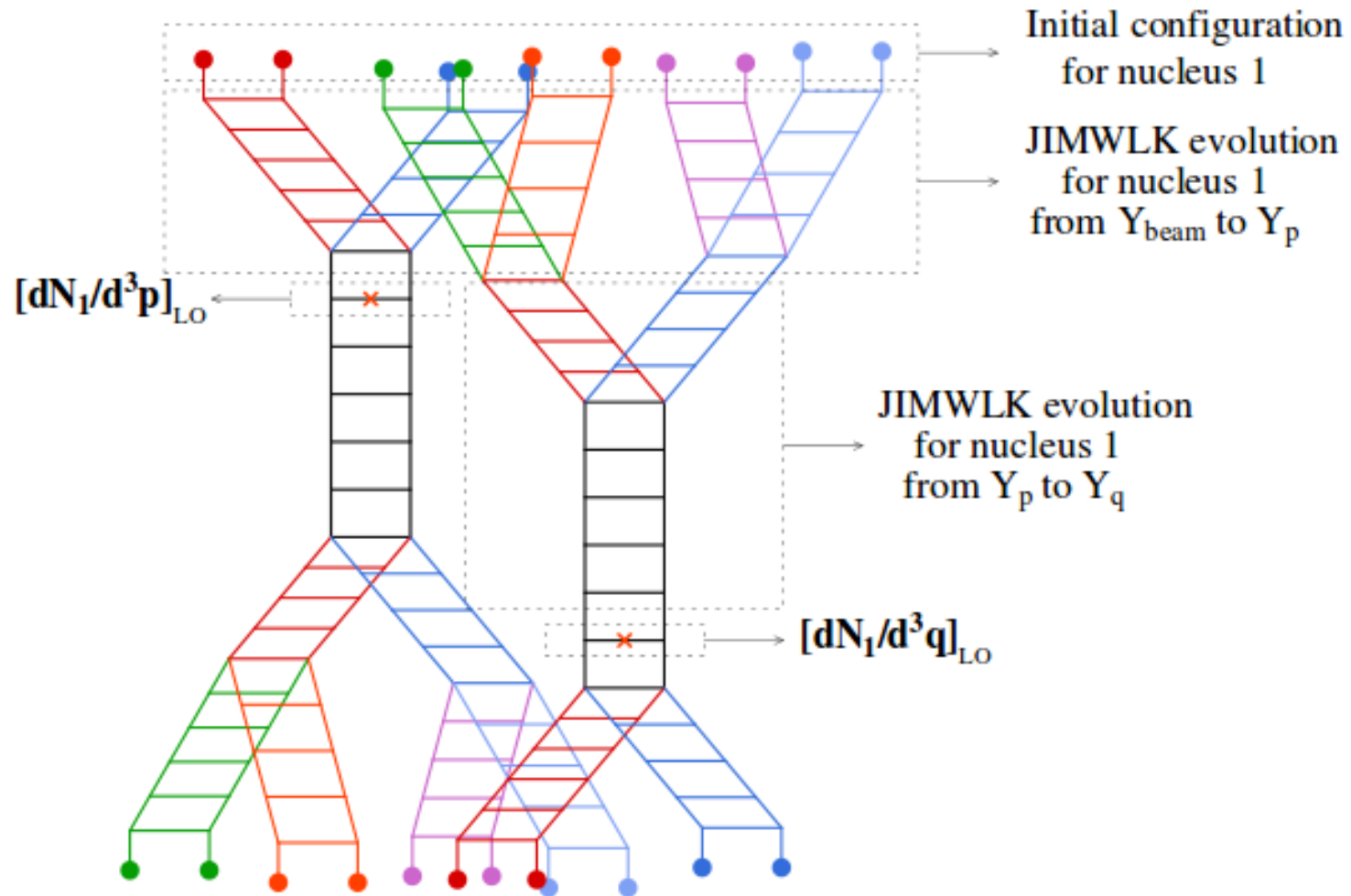
Forward jet structure



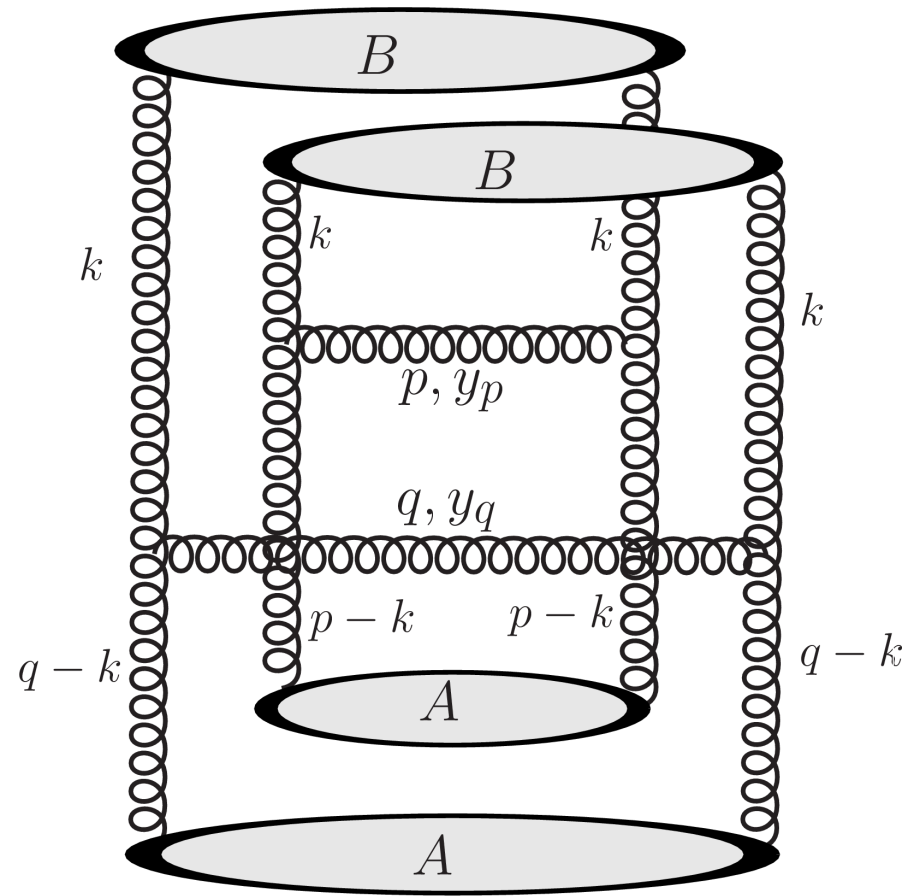
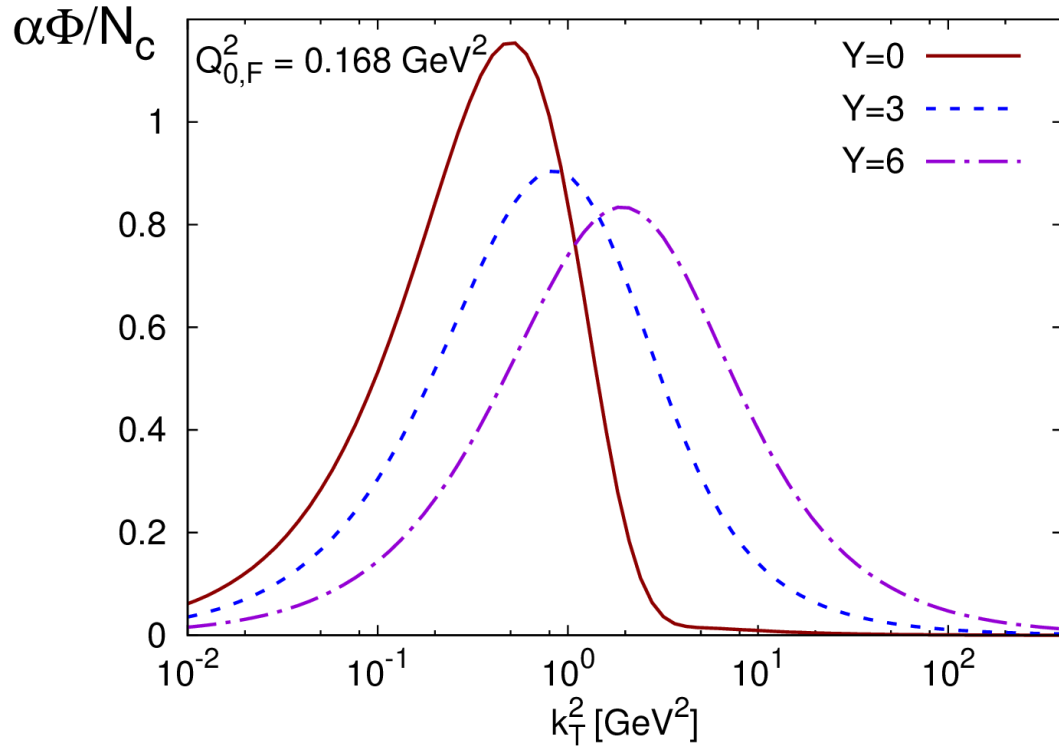
k_T factorization: double gluon production

$$\begin{aligned}
 & \left\langle \frac{dN_2}{d^2\mathbf{p}_\perp dy_p d^2\mathbf{q}_\perp dy_q} \right\rangle_{\text{LLog}} = \frac{32\alpha_s(\mathbf{p}_\perp)\alpha_s(\mathbf{q}_\perp)}{(2\pi)^{10} N_c C_F^3 \zeta} \frac{1}{\mathbf{p}_\perp^2 \mathbf{q}_\perp^2} \\
 & \times \left\{ \int d^2\mathbf{k}_{1\perp} \Phi_{A_1}^2(y_p, \mathbf{k}_{1\perp}) \Phi_{A_2}(y_p, \mathbf{p}_\perp - \mathbf{k}_{1\perp}) \Phi_{A_2}(y_q, \mathbf{q}_\perp + \mathbf{k}_{1\perp}) \right. \\
 & + \int d^2\mathbf{k}_{1\perp} \Phi_{A_1}^2(y_p, \mathbf{k}_{1\perp}) \Phi_{A_2}(y_p, \mathbf{p}_\perp - \mathbf{k}_{1\perp}) \Phi_{A_2}(y_q, \mathbf{q}_\perp - \mathbf{k}_{1\perp}) \\
 & + \int d^2\mathbf{k}_{1\perp} \Phi_{A_2}^2(y_q, \mathbf{k}_{1\perp}) \Phi_{A_1}(y_p, \mathbf{p}_\perp - \mathbf{k}_{1\perp}) \Phi_{A_1}(y_q, \mathbf{q}_\perp + \mathbf{k}_{1\perp}) \\
 & \left. + \int d^2\mathbf{k}_{1\perp} \Phi_{A_2}^2(y_q, \mathbf{k}_{1\perp}) \Phi_{A_1}(y_p, \mathbf{p}_\perp - \mathbf{k}_{1\perp}) \Phi_{A_1}(y_q, \mathbf{q}_\perp - \mathbf{k}_{1\perp}) \right\}
 \end{aligned}$$

k_T factorization: double gluon production



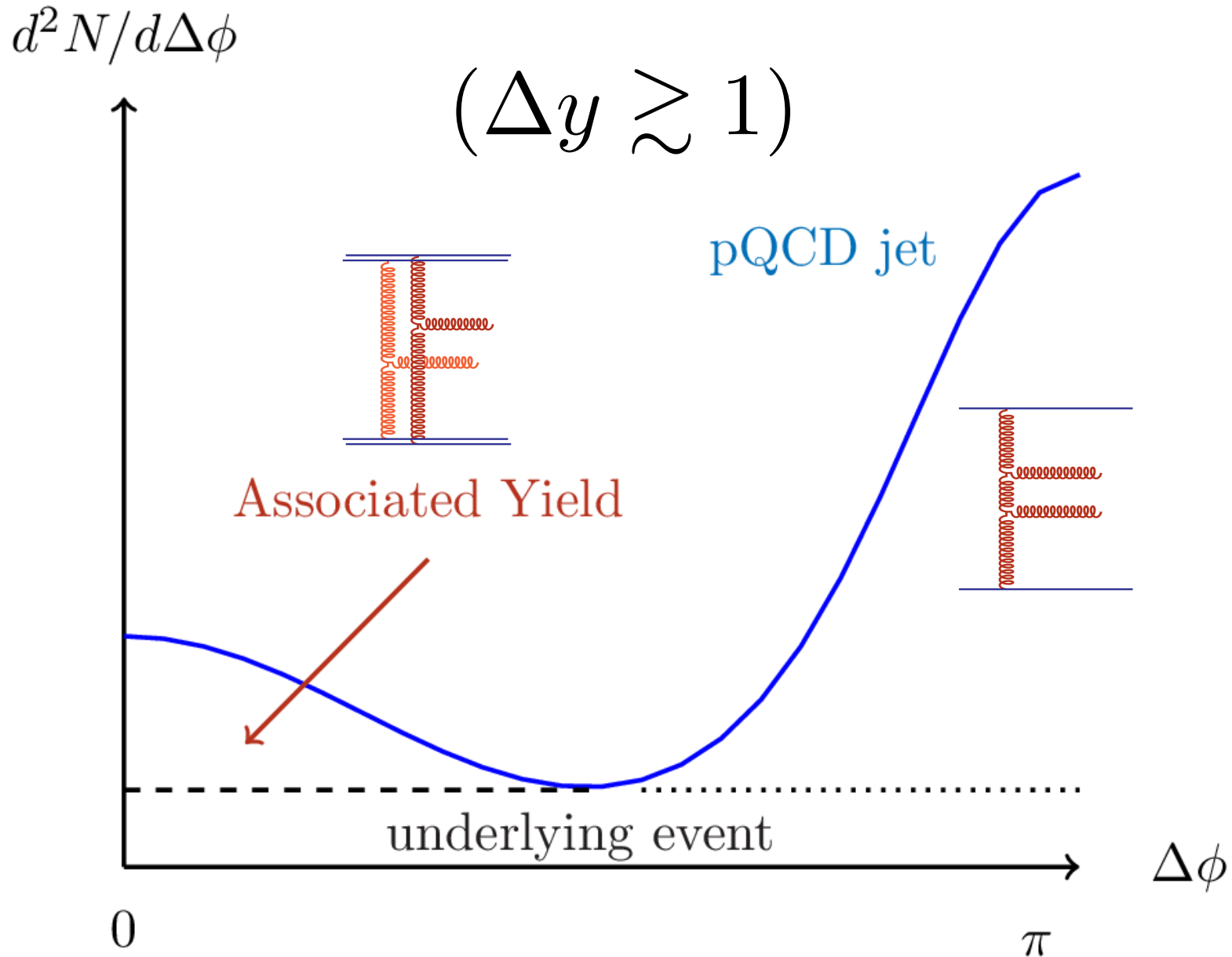
Angular Structure



Condition for Ridge (Qualitatively):

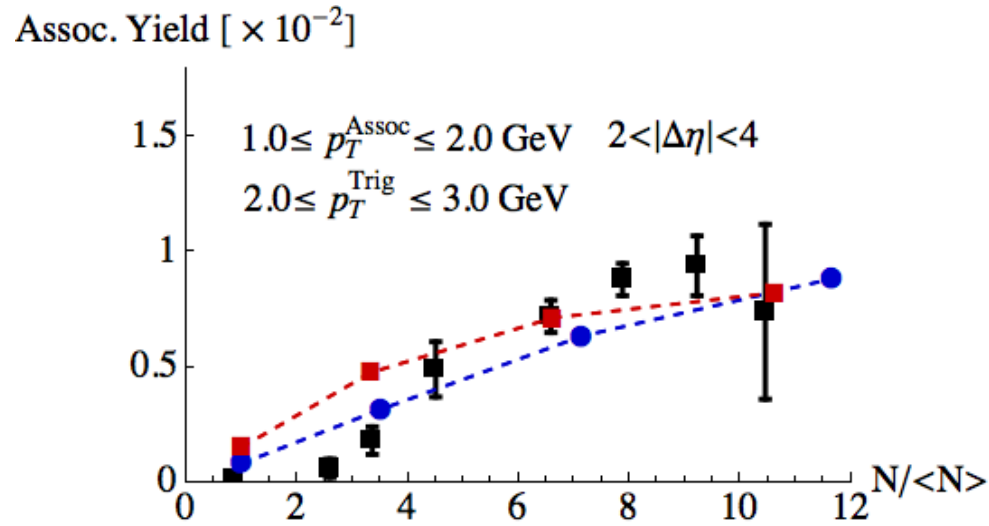
$$|\mathbf{k}_{\perp}| \sim |\mathbf{p}_{\perp} - \mathbf{k}_{\perp}| \sim |\mathbf{q}_{\perp} \pm \mathbf{k}_{\perp}| \sim Q_s$$

Forward jet structure



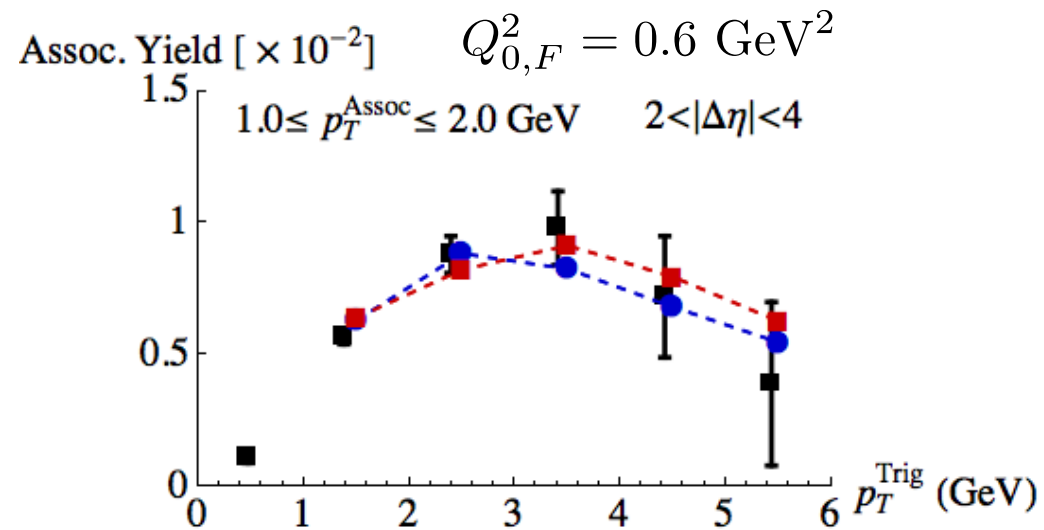
Ridge in p+p collisions

Centrality Dependence:

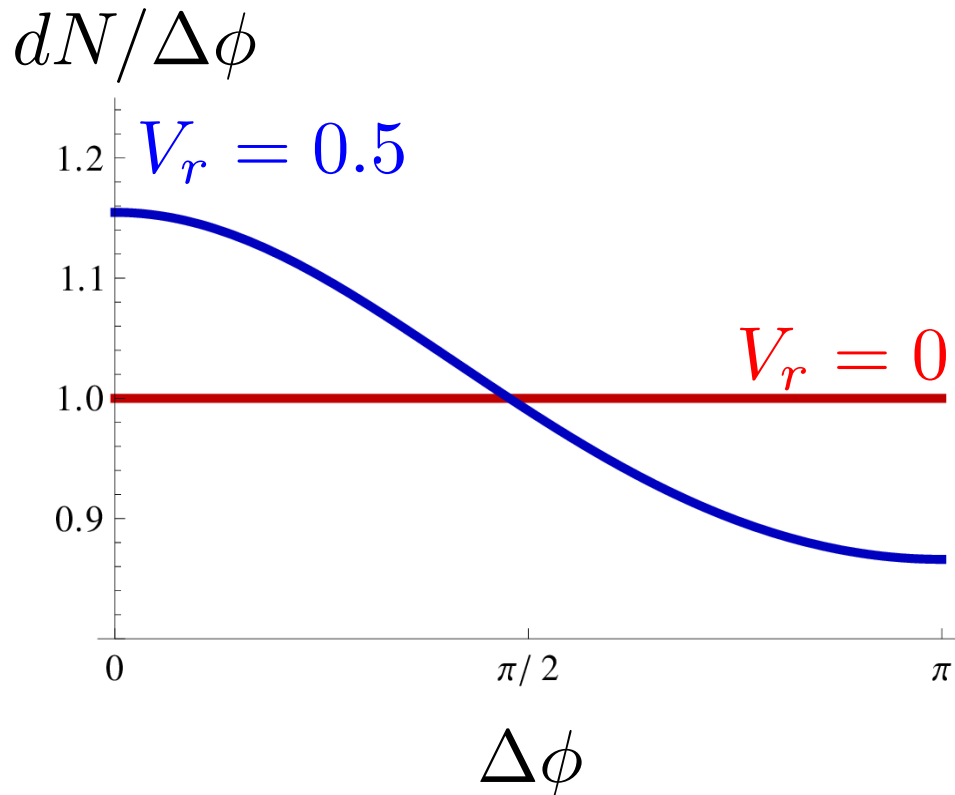


$$Q_{0,F}^2(x_0 = 0.01) = 0.15, 0.3, 0.45, 0.6 \text{ GeV}^2$$

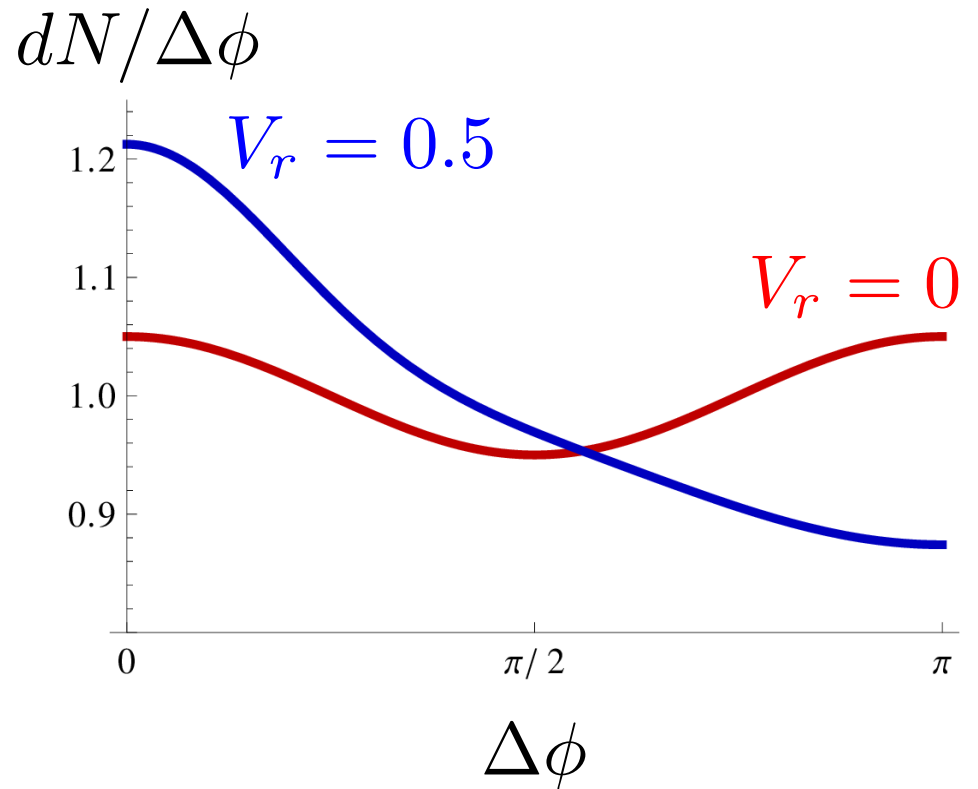
Trigger Dependence:



Role played by transverse flow



Left: No intrinsic correlation in $\Delta\phi$ followed by radial boost.



Right: Intrinsic azimuthal correlation followed by boost.

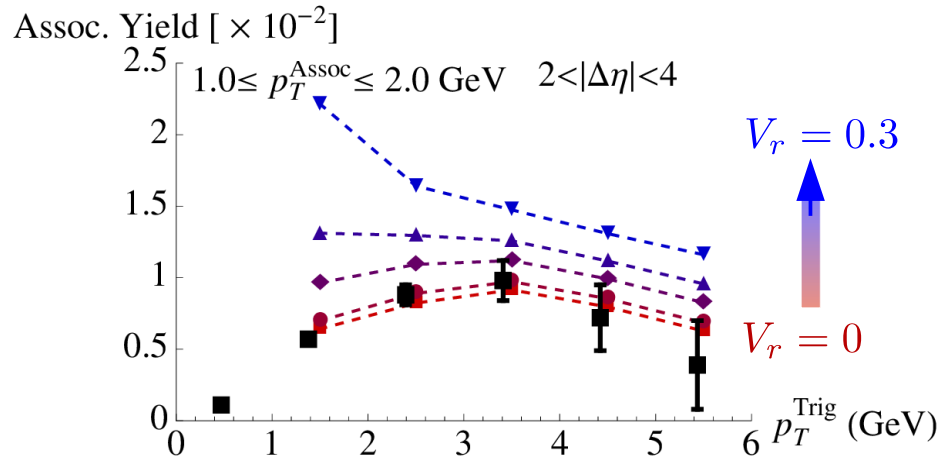
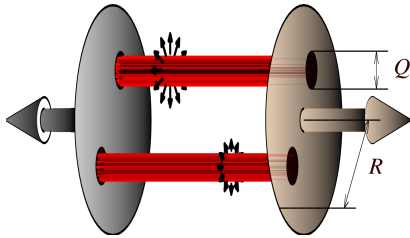
Transverse flow increases near side collimation;
but is it seen in the data?

p+p

vs

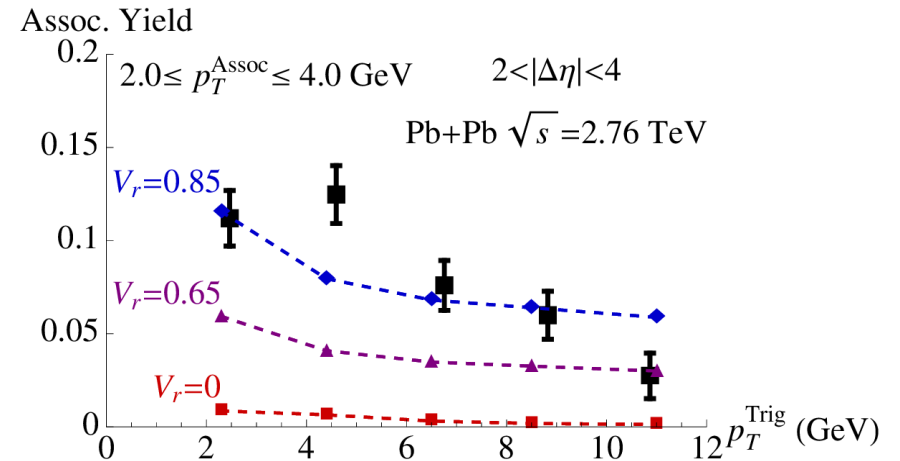
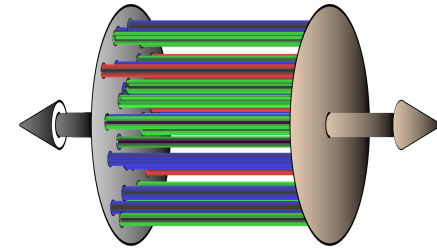
A+A

In p+p we are seeing the intrinsic collimation from a single flux tube



Increasing transverse flow in p+p creates a discrepancy with data.

In A+A there are many such tubes each with an intrinsic correlation enhanced by flow

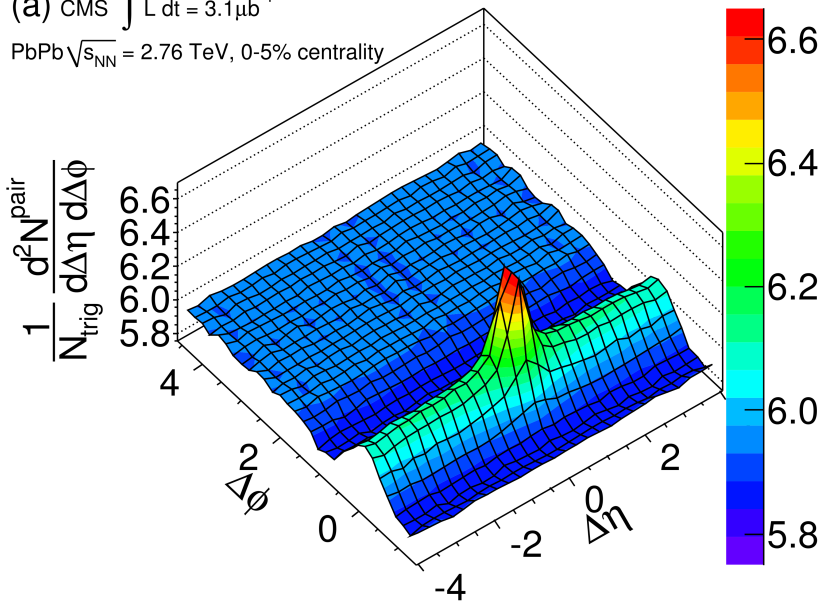


Yet, transverse flow is needed to explain identical measurements in Pb+Pb

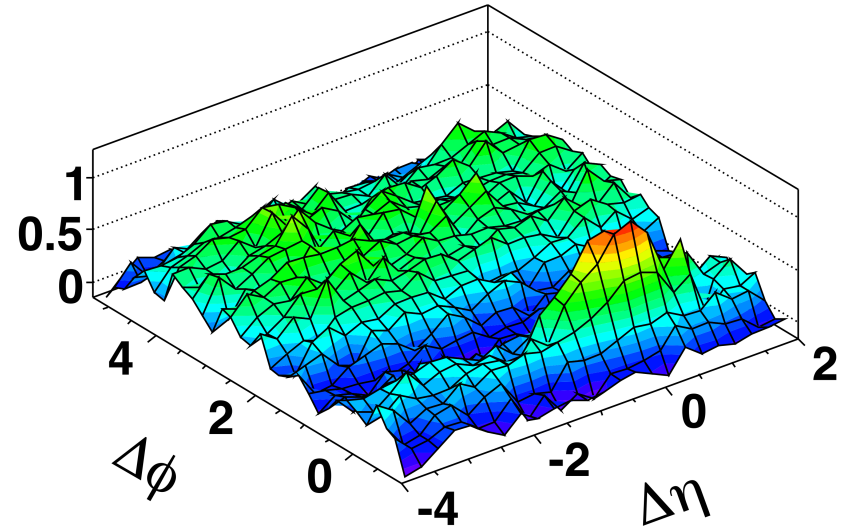
Are we sure the A+A ridge is probing the nuclear wavefunction?

Heavy-Ion Ridge

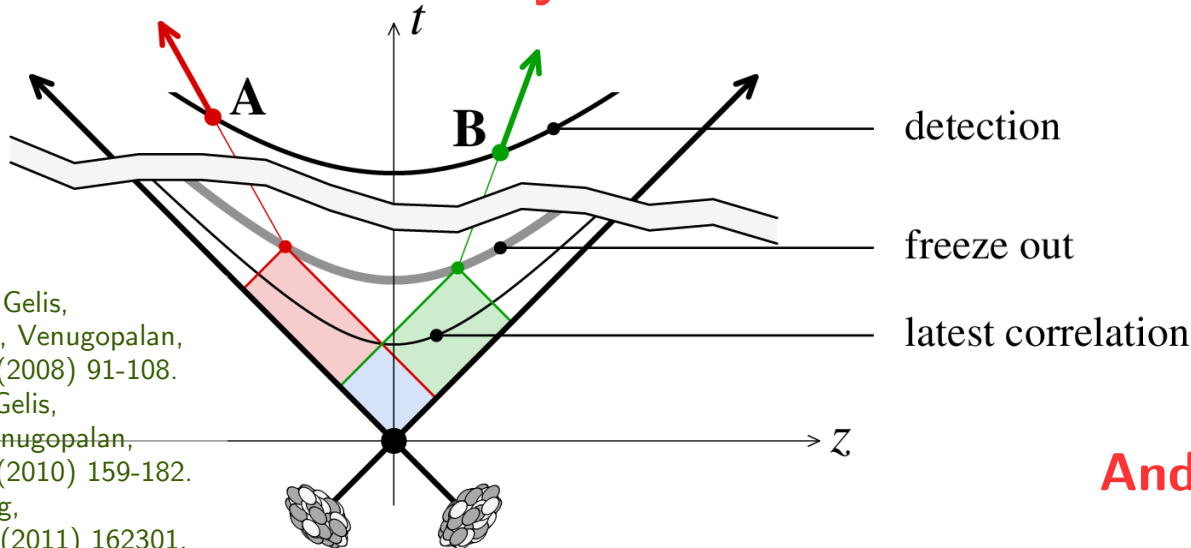
(a) CMS $\int L dt = 3.1 \mu\text{b}^{-1}$
 PbPb $\sqrt{s_{NN}} = 2.76 \text{ TeV}$, 0-5% centrality



PHOBOS (Au+Au) $\sqrt{s} = 200 \text{ GeV}$



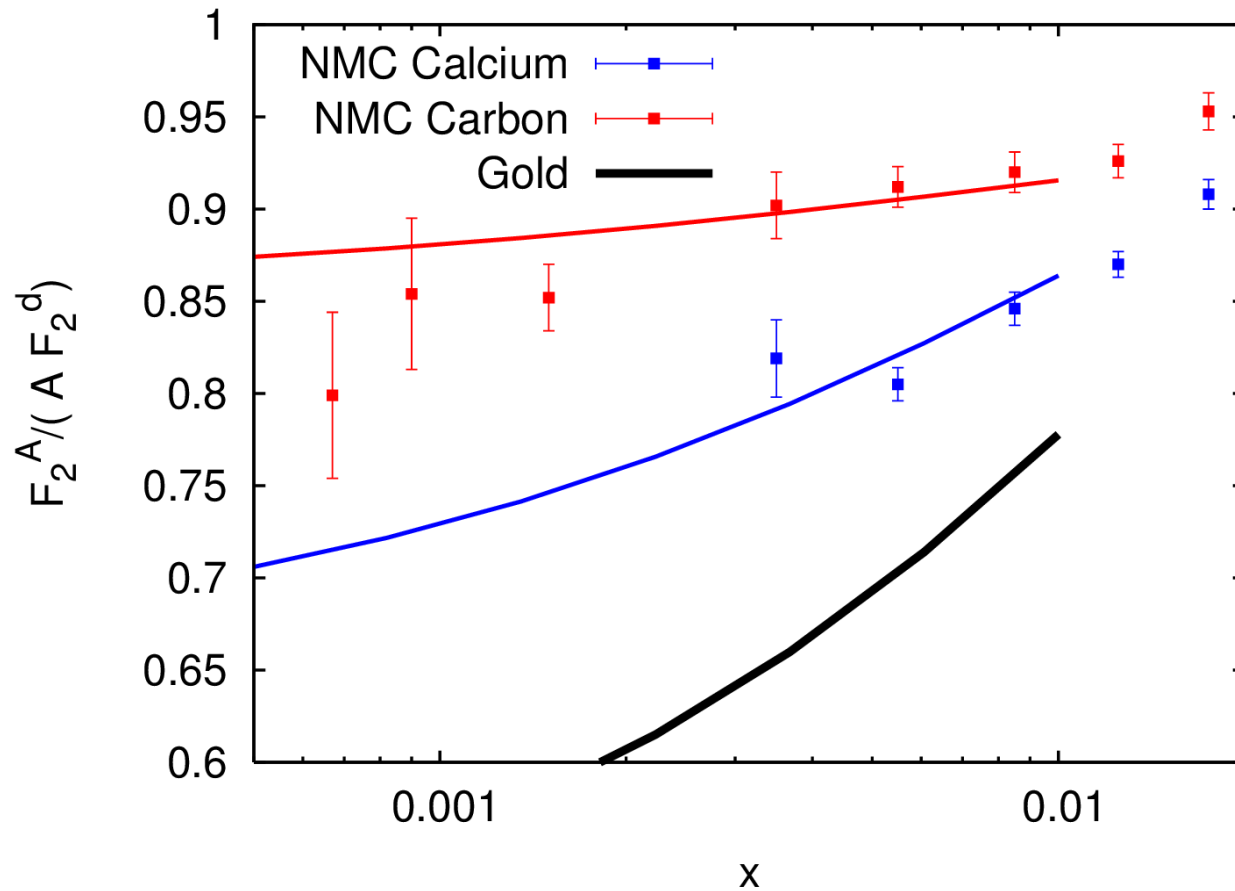
The correlation is long range in rapidity.
 Causality dictates the correlation formed early.



Dumitru, Gelis,
 McLerran, Venugopalan,
 NPA810 (2008) 91-108.
 Dusling, Gelis,
 Lappi, Venugopalan,
 NPA836 (2010) 159-182.
 Ma, Wang,
 PRL 106 (2011) 162301.

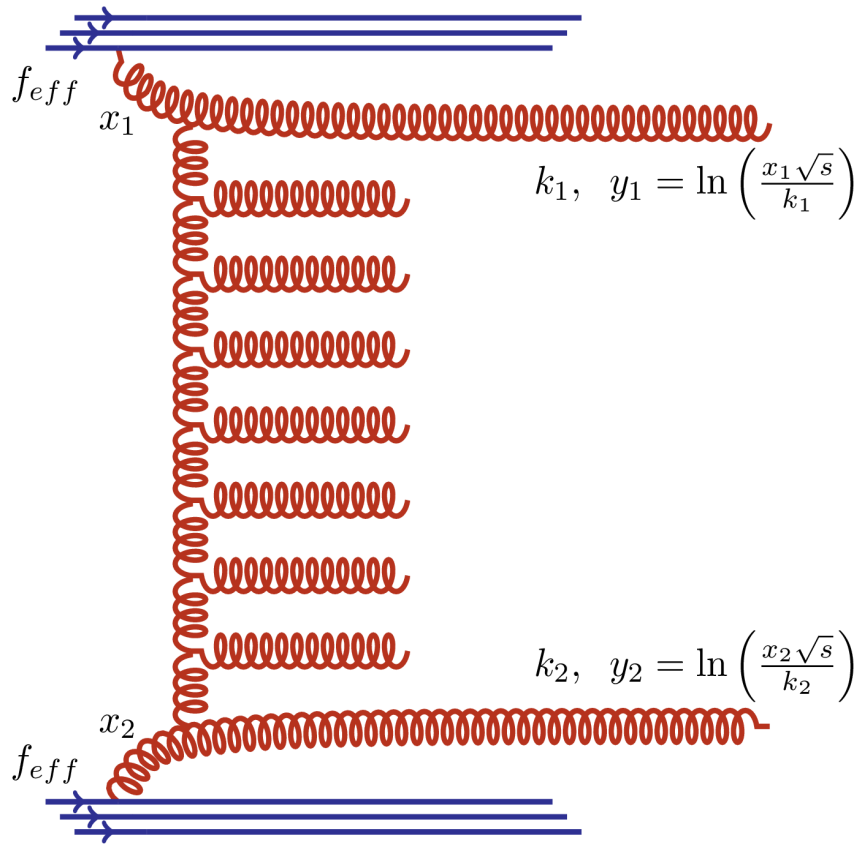
And it persists to the final state:

Need for EIC:



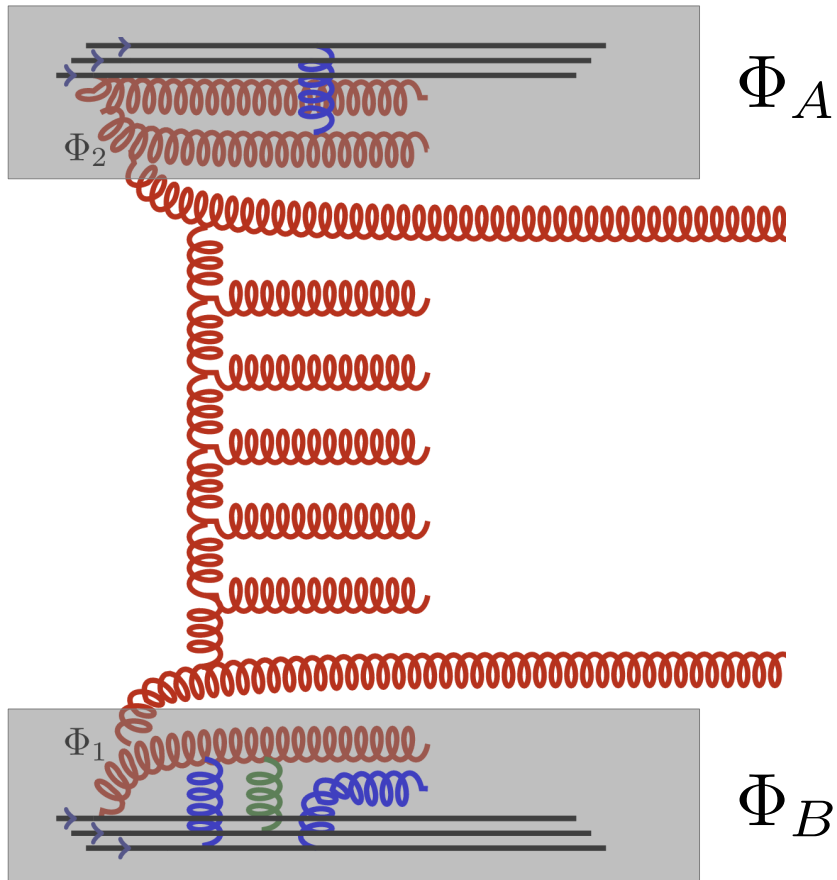
World collection of small x ($x \leq 0.01$) data for DIS on nuclei.

Part II: Jet Decorrelation



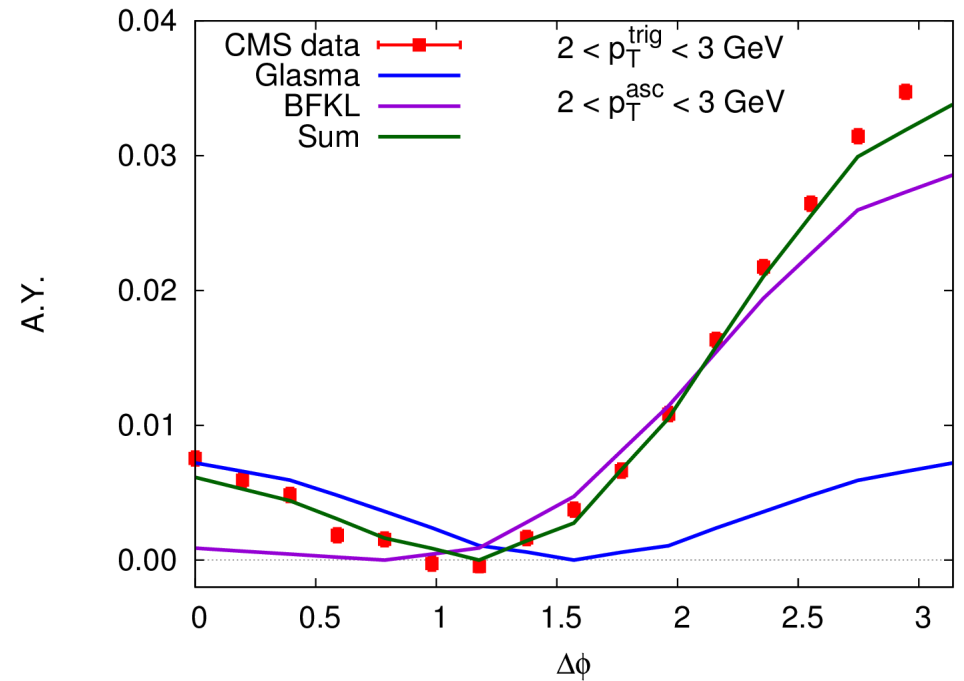
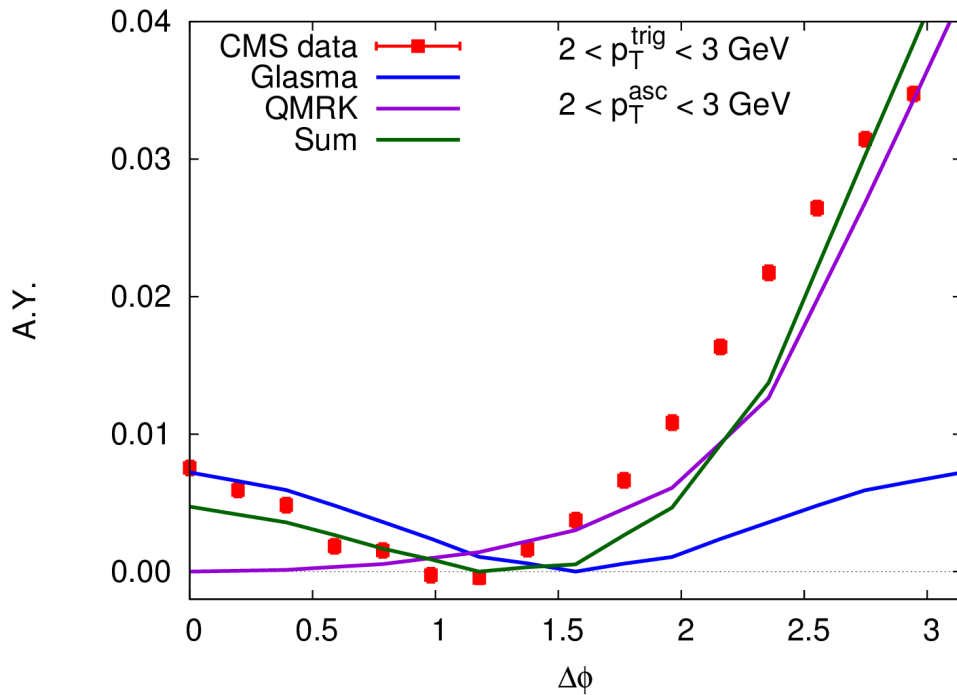
One can look for the growth in cross section with larger rapidity gaps as first suggested by Muller and Navelet.

Mini jets in k_T factorization

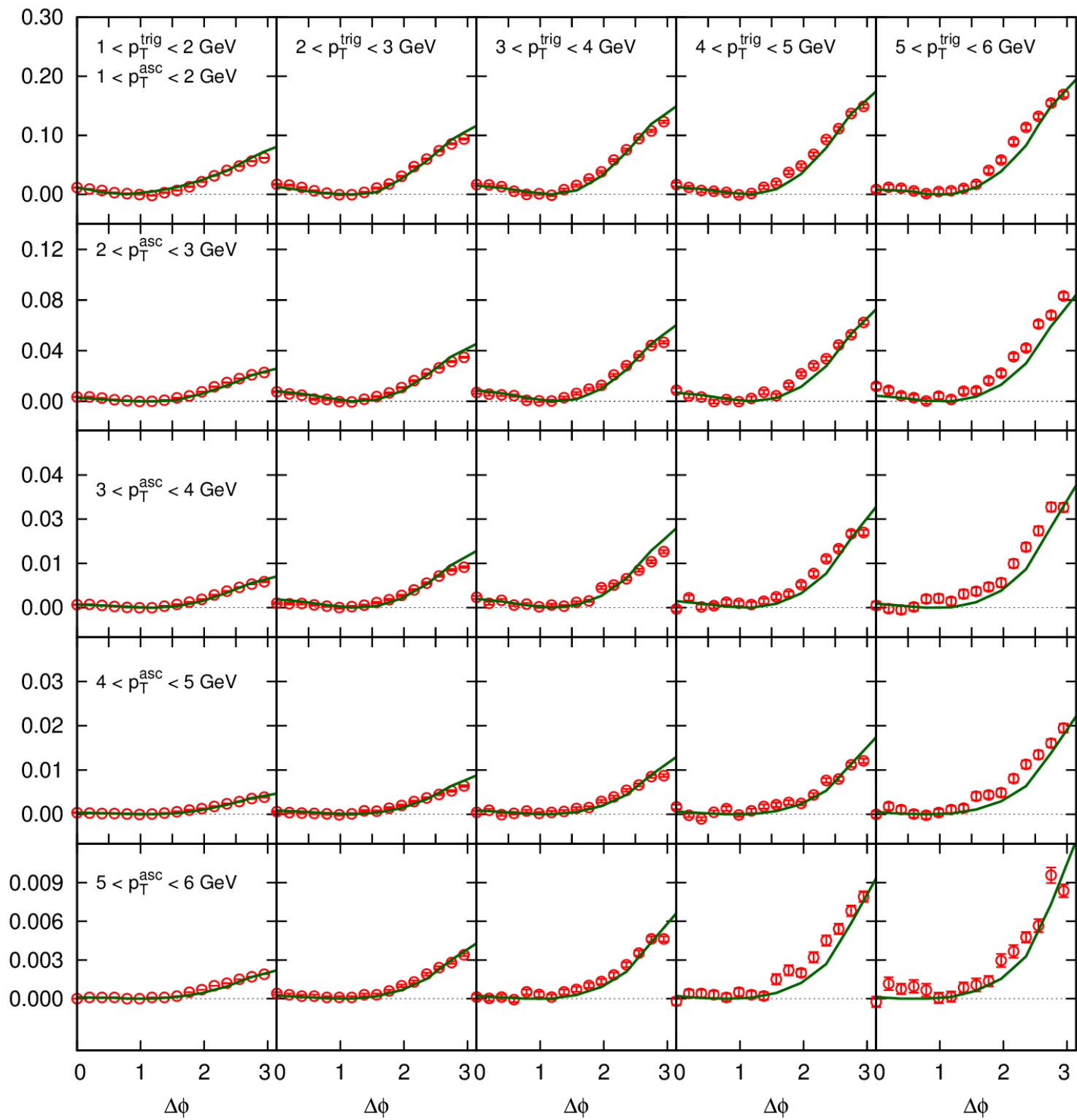


$$\frac{d^2 N_{AB}}{d^2 \mathbf{p}_\perp d^2 \mathbf{q}_\perp dy_p dy_q} = \frac{32 N_c \alpha_s(\mathbf{p}_\perp) \alpha_s(\mathbf{q}_\perp)}{(2\pi)^8 C_F} \frac{1}{\mathbf{p}_\perp^2 \mathbf{q}_\perp^2} \times \int d^2 \mathbf{k}_{0\perp} \int d^2 \mathbf{k}_{3\perp} \Phi_A(x_1, \mathbf{k}_{0\perp}) \Phi_B(x_2, \mathbf{k}_{3\perp}) \mathcal{G}(\mathbf{k}_{0\perp} - \mathbf{p}_\perp, \mathbf{k}_{3\perp} + \mathbf{q}_\perp, y_p - y_q)$$

Evidence for BFKL evolution

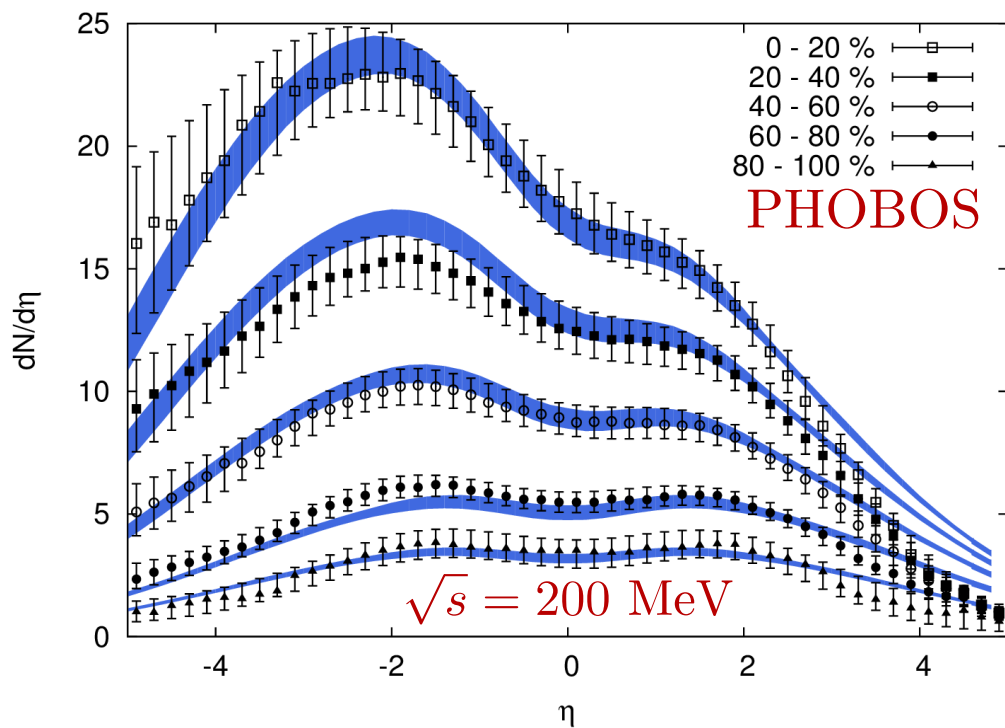


There is a clear need for evolution between the triggered particles (even for a rapidity gap as small as 2-4 units)

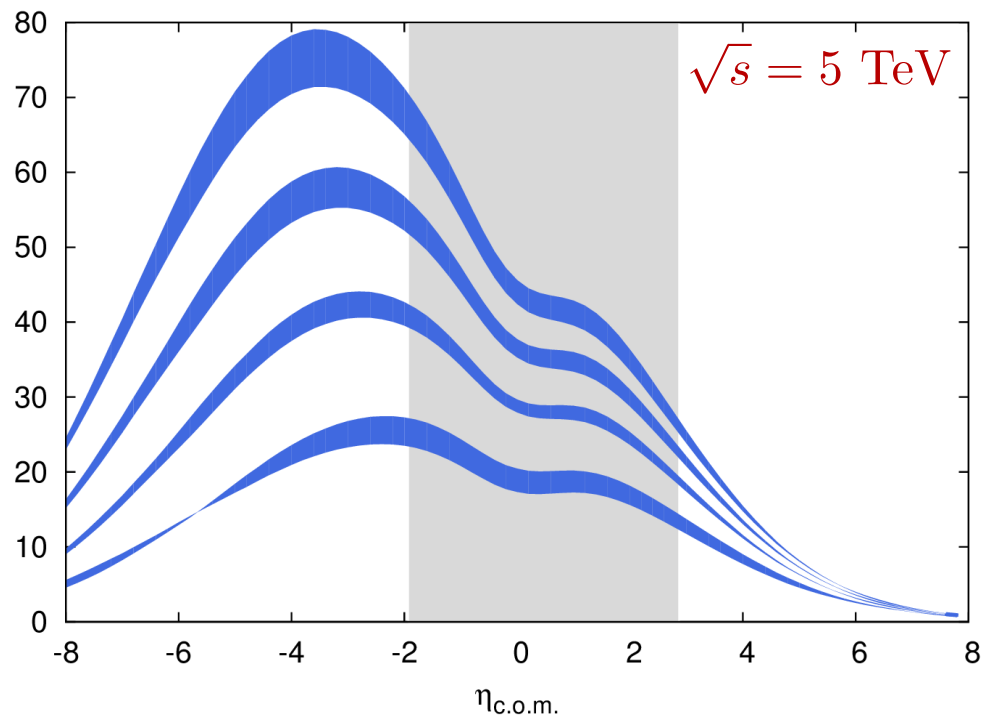


Part 3: Prospects for p+Pb

$$Q_{0,Au}^2 = 0.336, 0.336, 0.504, 0.672, 0.840 \text{ GeV}^2 \quad Q_{0,Pb}^2 = 0.504, 0.672, 0.840, 1.008 \text{ GeV}^2$$

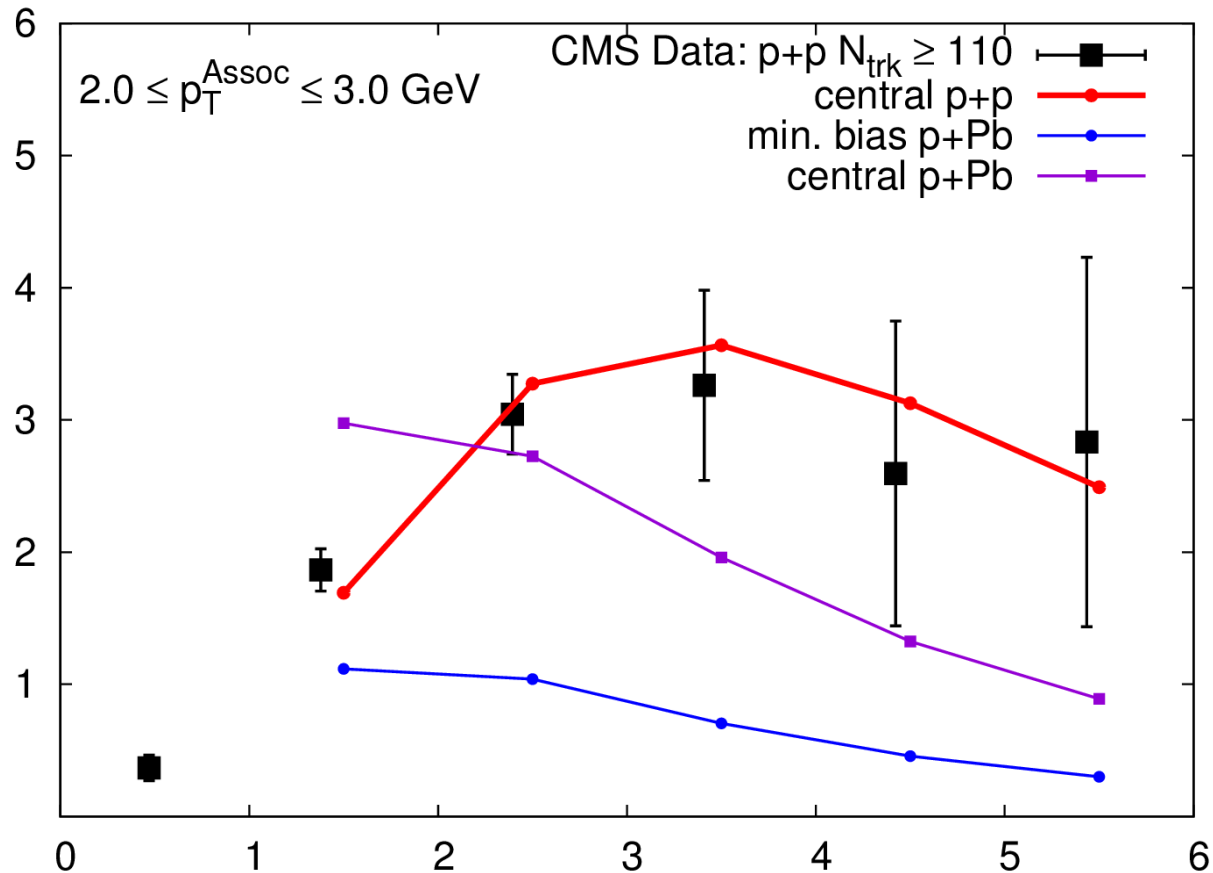


$$Q_{0,d}^2 = 0.336 \text{ GeV}^2$$



$$Q_{0,p}^2 = 0.168 \text{ GeV}^2$$

Ridge in p+Pb



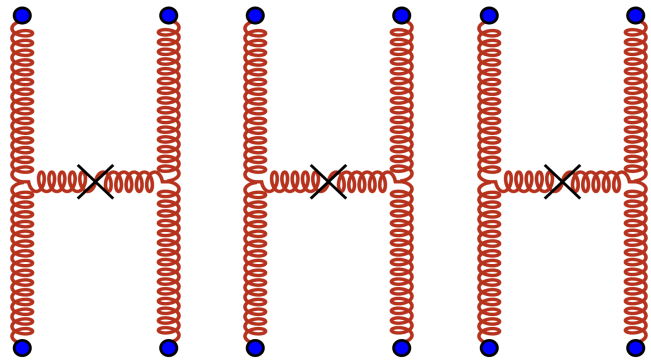
Ridge in p+Pb is smaller than in p+p for CMS acceptance. Signal will also have to be pulled from a larger background.

Summary

- Strong color sources lead to α_s^8 enhancement of QCD diagram responsible for near-side enhancement
- Near side collimation is a clear signature of saturation dynamics
- Clear evidence for BFKL evolution in CMS dijet measurements

Backup

High multiplicity are b=0 collisions



$$P_n^{\text{NB}}(\bar{n}, k) = \frac{\Gamma(k+n)}{\Gamma(k)\Gamma(n+1)} \frac{\bar{n}^n k^k}{(\bar{n}+k)^{n+k}}$$

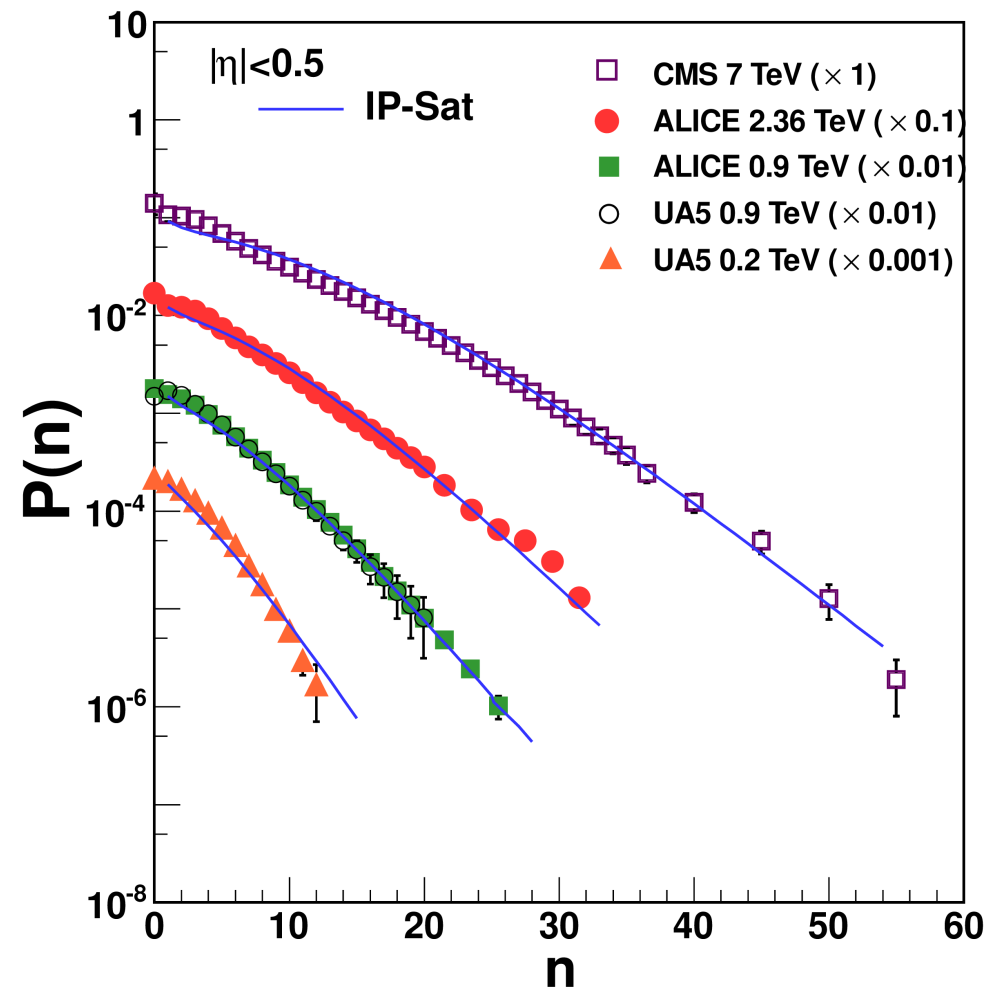
Dumitru, Gelis, McLerran, Venugopalan, NPA810 91-108 (2008).
 Dusling, Fernandez-Fraile, Venugopalan NPA828 (2009) 161-177.
 Gelis, Lappi, McLerran, NPA828 (2009) 149-160.

$$k = \zeta \frac{(N_c^2 - 1) S_{\perp} Q_s^2}{2\pi}$$

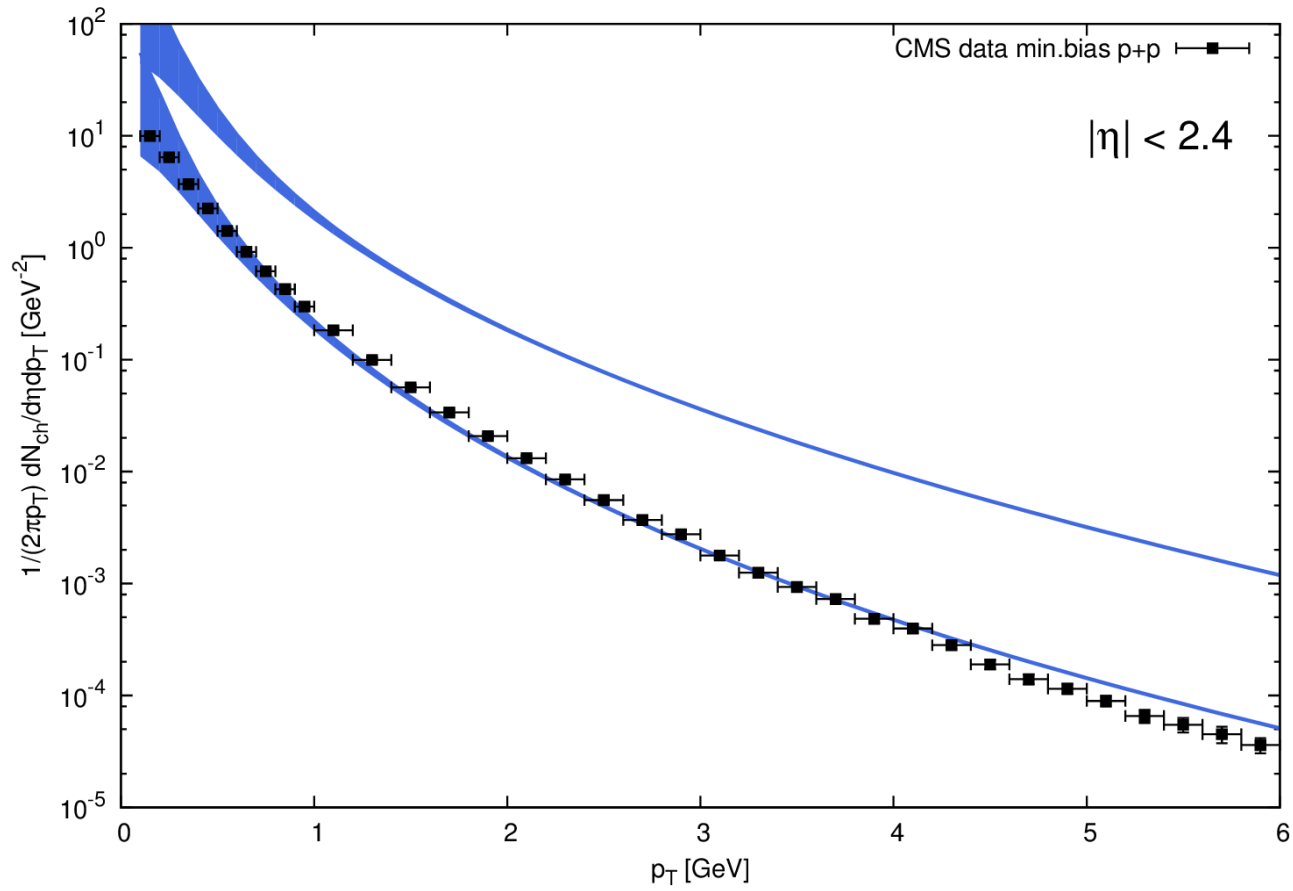
$$\zeta = 0.155 \text{ [Empirical]}$$

$$\zeta = 0.2 - 1.5 \text{ [Lattice]}$$

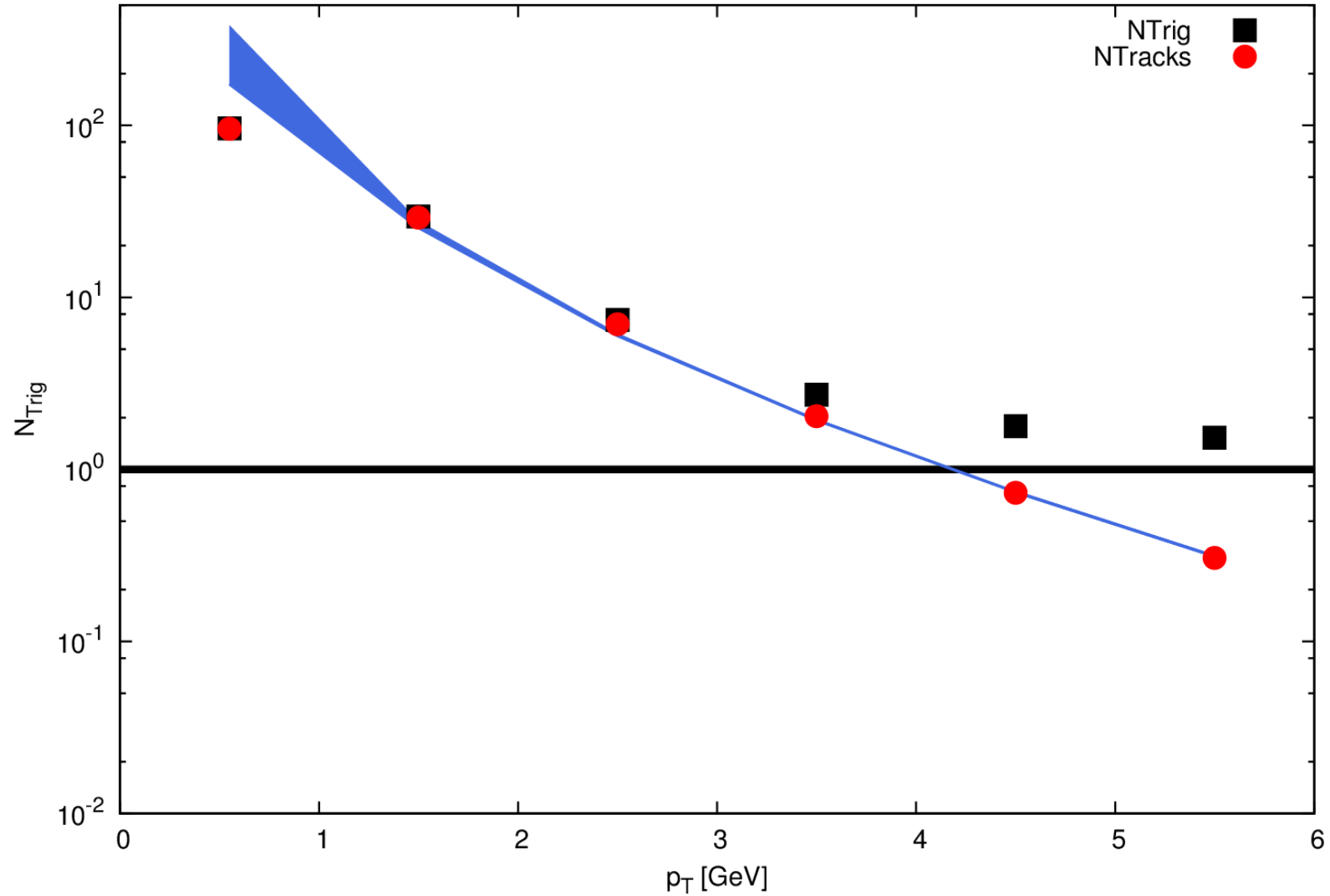
Empirical: Tribedy, Venugopalan, NPA850 (2011) 136-156.
 Lattice (CYM): Lappi, Srednyak, Venugopalan, JHEP01 (2010) 066.
 Schenke, Tribedy, Venugopalan, arXiv:1206.6805



$p+p$ p_T distribution



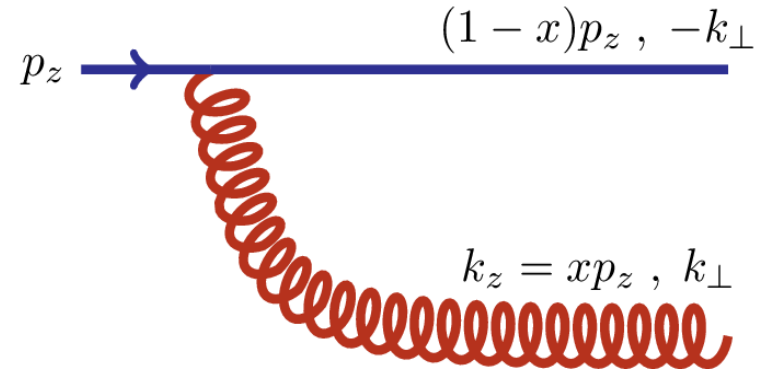
CMS Acceptance



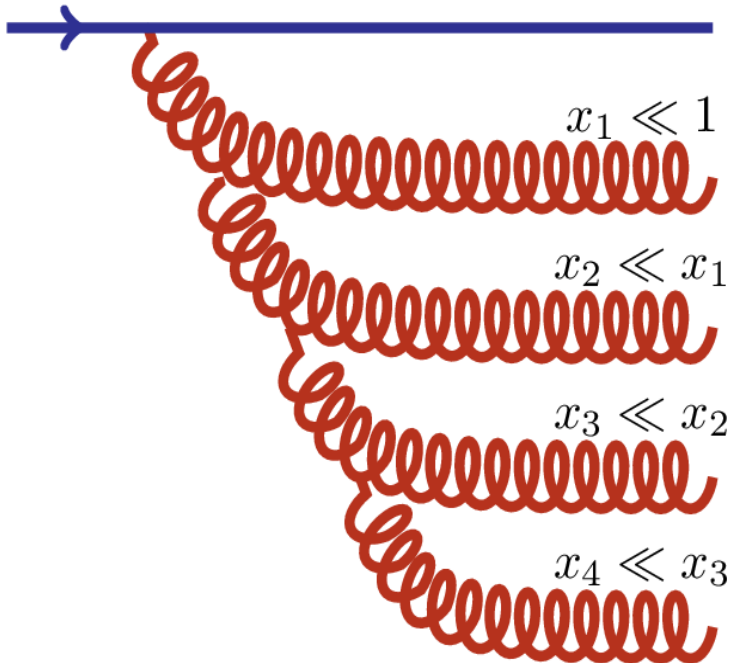
Gluon radiation

As the energy is increased new gluons are emitted with probability

$$dP_{\text{Brem}} \sim C_R \frac{\alpha_s}{\pi^2} \frac{d^2 k_{\perp}}{k_{\perp}^2} \frac{dx}{x}$$



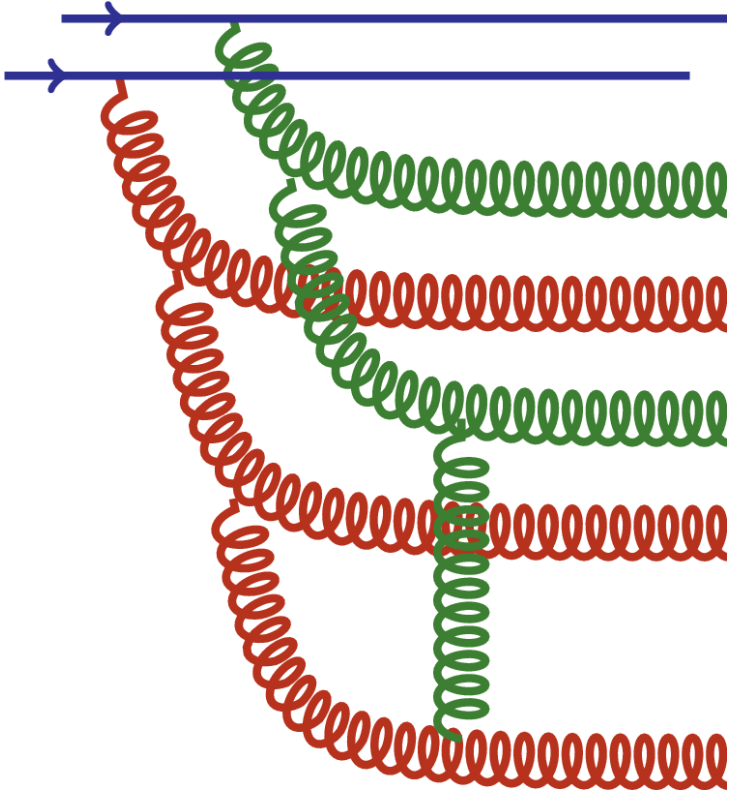
And as long as the density remains low the evolution is linear



$$\frac{\partial T(\mathbf{r}_{\perp}, Y)}{\partial Y} = \frac{\alpha_s N_c}{2\pi^2} \int d\mathbf{r}_1 \frac{\mathbf{r}_{\perp}^2}{\mathbf{r}_1^2 (\mathbf{r}_{\perp} - \mathbf{r}_1)^2} \times [T(\mathbf{r}_1, Y) + T(\mathbf{r}_{\perp} - \mathbf{r}_1, Y) - T(\mathbf{r}_{\perp}, Y)]$$

Kuraev, Lipatov, Fadin, Sov.Phys.JETP44 443-450 (1976).
Sov.Phys.JETP45 199-204 (1977).
Balitsky, Lipatov, Sov.J.Nucl.Phys 28 822-829 (1978).

BK Evolution Equation



Balitsky, NPB 463, 99 (1996).
Kovchegov, PRD 60, 034008 (1999).

Jalilian-Marian, Kovner, McLerran, Weigert, PRD 55 5414 (1997).
Jalilian-Marian, Kovner, Leonidov, Weigert, NPB 504 415 (1997),
PRD 59 014014 (1999).

$$\frac{\partial T(\mathbf{r}_\perp, Y)}{\partial Y} = \frac{\alpha_s N_c}{2\pi^2} \int d\mathbf{r}_1 \frac{\mathbf{r}_\perp^2}{\mathbf{r}_1^2 (\mathbf{r}_\perp - \mathbf{r}_1)^2} \times$$
$$[T(\mathbf{r}_1, Y) + T(\mathbf{r}_\perp - \mathbf{r}_1, Y) - T(\mathbf{r}_\perp, Y) - T(\mathbf{r}_1, Y)T(\mathbf{r}_\perp - \mathbf{r}_1, Y)]$$

NLO BK Equation

$$\frac{\partial T(\mathbf{r}_\perp, Y)}{\partial Y} = \int d\mathbf{r}_1 \mathcal{K}_{\text{Bal.}}(\mathbf{r}_\perp, \mathbf{r}_1, \mathbf{r}_\perp - \mathbf{r}_1) \times \\ [T(\mathbf{r}_1, Y) + T(\mathbf{r}_\perp - \mathbf{r}_1, Y) - T(\mathbf{r}_\perp, Y) - T(\mathbf{r}_1, Y)T(\mathbf{r}_\perp - \mathbf{r}_1, Y)]$$

$$\mathcal{K}_{\text{Bal.}}(\mathbf{r}, \mathbf{r}_1, \mathbf{r}_2) = \frac{\alpha_s(\mathbf{r}) N_c}{\pi} \left[\frac{\mathbf{r}^2}{\mathbf{r}_1^2 \mathbf{r}_2^2} + \frac{1}{\mathbf{r}_1^2} \left(\frac{\alpha_s(\mathbf{r}_1^2)}{\alpha_s(\mathbf{r}_2^2)} - 1 \right) + \frac{1}{\mathbf{r}_2^2} \left(\frac{\alpha_s(\mathbf{r}_2^2)}{\alpha_s(\mathbf{r}_1^2)} - 1 \right) \right]$$

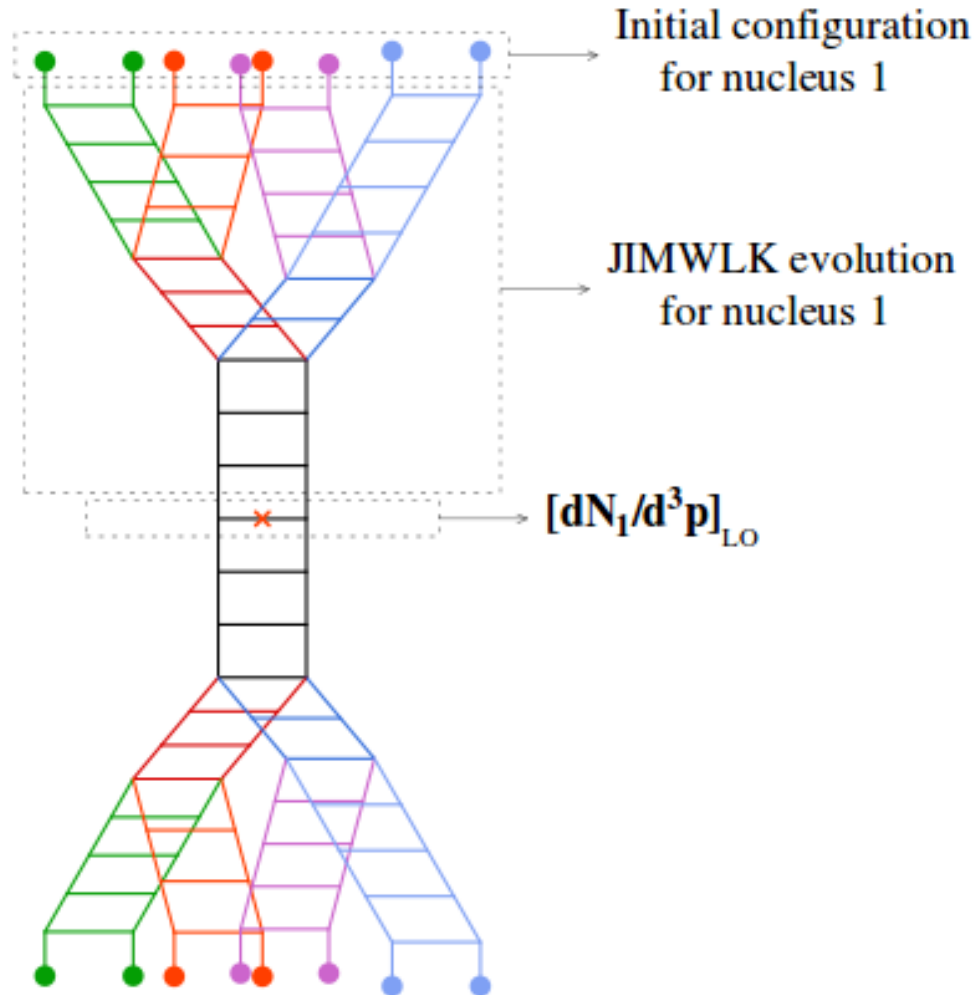
Blast Wave Model

$$\frac{d^2 N}{d\Delta\phi} = \int_{-\pi}^{\pi} d\Psi \mathcal{J}(\Psi, \Delta\phi) \frac{d^2 N}{d\Delta\tilde{\phi}}(\Delta\tilde{\phi}(\Psi, \Delta\phi))$$

$$2 \sin^2 \left(\frac{\Delta\tilde{\phi}}{2} \right) = \frac{\sqrt{1 - \beta^2} (1 - \cos(\Delta\phi))}{1 - 2\beta \cos \Psi \cos \left(\frac{\Delta\phi}{2} \right) + \frac{\beta^2}{2} (\cos(\Delta\phi) + \cos(2\Psi))} .$$

$$\mathcal{J} = \frac{1 - \beta^2}{(1 - \beta \cos(\Psi + \Delta\phi/2))(1 - \beta \cos(\Psi - \Delta\phi/2))} ,$$

k_T factorization: single gluon production



$$\left\langle \frac{dN_1}{d^2\mathbf{p}_\perp dy_p} \right\rangle_{LLog} = \frac{8\alpha_s(p_\perp)S_\perp}{C_F(2\pi)^4} \frac{1}{\mathbf{p}_\perp^2} \int \frac{d^2k_\perp}{(2\pi)^2} \Phi_{A_1}(y_p, k_\perp) \Phi_{A_2}(y_p, p_\perp - k_\perp)$$

BFKL Formalism

$$\mathcal{G}(a, b, \Delta y) = \frac{1}{(2\pi)^2} \frac{1}{(a^2 b^2)^{1/2}} \sum_n e^{in\bar{\phi}} \int_{-\infty}^{+\infty} d\nu e^{\omega(\nu, n)\Delta y} e^{i\nu \ln(a^2/b^2)}$$

$$\omega(\nu, n) = -2\bar{\alpha}_s \operatorname{Re} \left[\Psi \left(\frac{|n|+1}{2} + i\nu \right) - \Psi(1) \right]$$

$$\bar{\alpha}_s \equiv \frac{N_c \alpha_s \left(\sqrt{ab} \right)}{\pi}$$

$$\bar{\phi} \equiv \arccos \left(\frac{a \cdot b}{|a| |b|} \right)$$